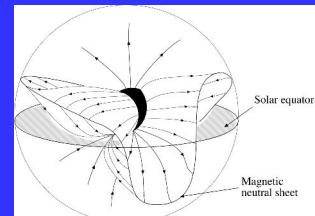
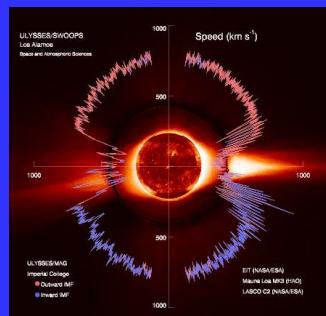


The solar corona and the stream structure of the solar wind

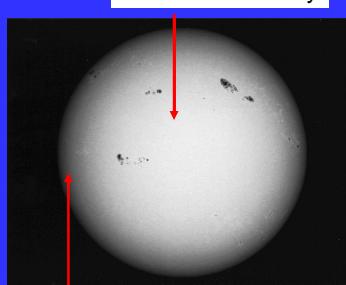
Hannu Koskinen

with special thanks to prof. Rainer Schwenn, MPS, Lindau



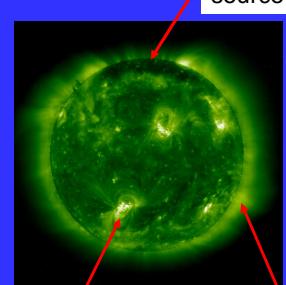
Solar atmosphere

Cool surface
5778 K black body



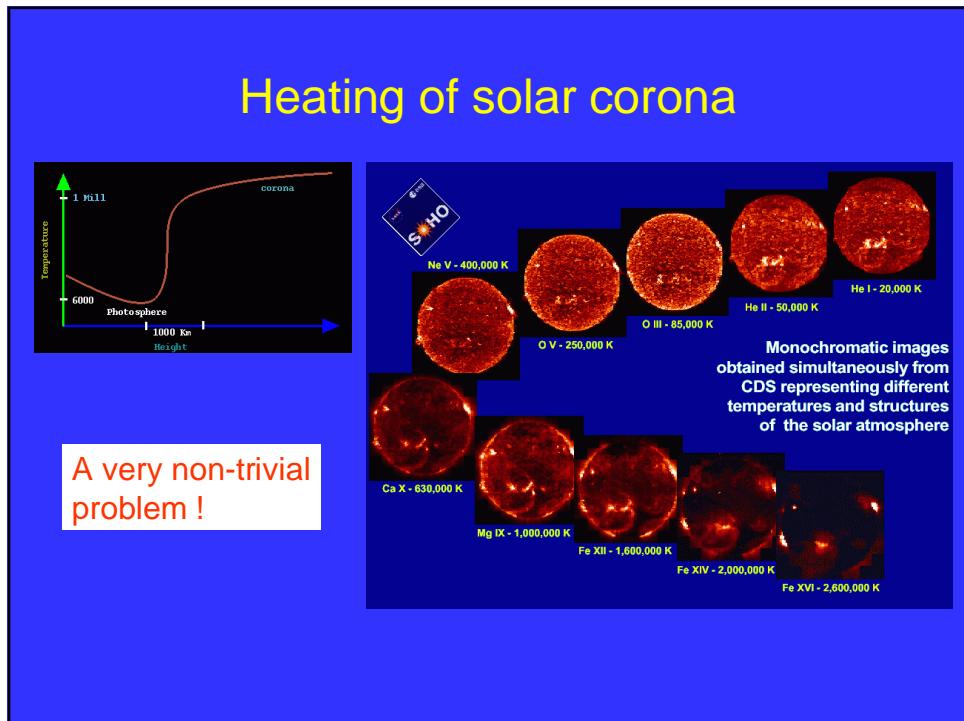
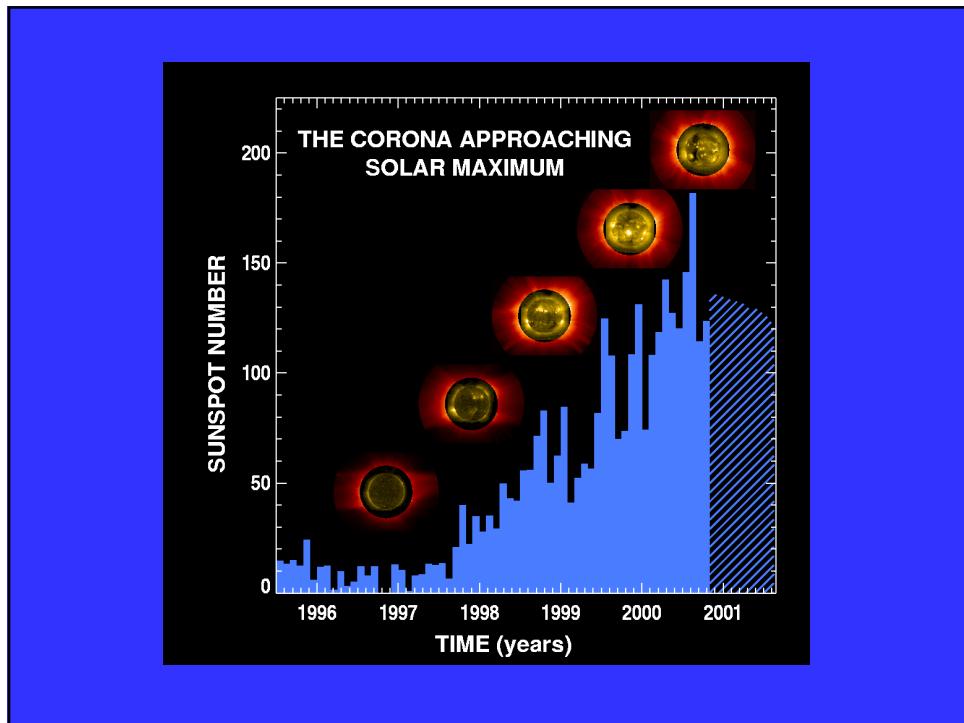
Limb darkening
Sun has an atmosphere

Coronal hole
source of solar wind



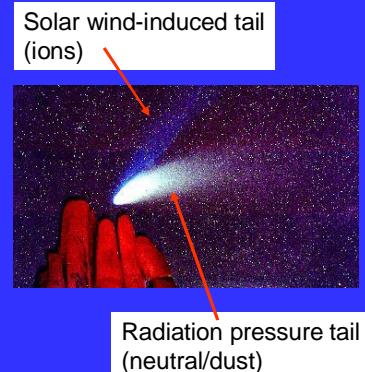
Various types
of activity

Hot corona
up to 2 MK



Why solar wind?

- There must be a mechanism to transfer solar activity to the Earth
 - Chapman 1929: Solar flares emit plasma clouds
 - Chapman and Ferraro ~ 1931: First ideas of magnetic storms
 - Biermann 1951: Cometary tails require a fast corpuscular flow in addition to radiation pressure



Parker's solution: Isothermal expanding solar wind

Assume a spherically symmetric time-independent outward flow.
Equations of continuity, momentum and state are:

$$4\pi r^2 n v = \text{const}$$

$$nmv \frac{dv}{dr} = -\frac{dp}{dr} - \frac{GM_{\odot} mn}{r^2}$$

$$p = nk_B T$$

Although the solar wind certainly cools when expanding, assume that the expansion is isothermal ($T = \text{const.}$)

$$\mathbf{P} \left(v - \frac{v_c^2}{v} \right) \frac{dv}{dr} = \frac{2v_c^2}{r} - \frac{GM_{\odot}}{r^2} \quad \text{where } v_c = \sqrt{k_B T / m}$$

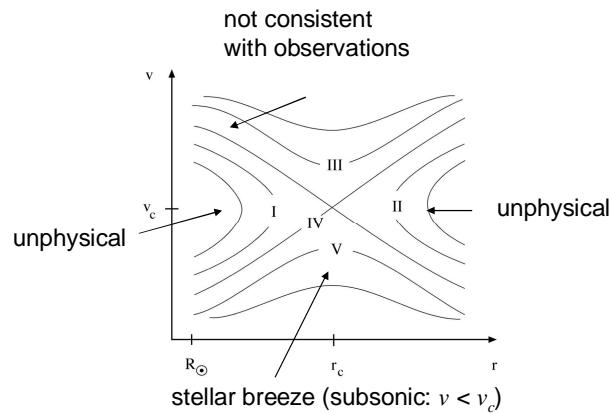
is the isothermal sound speed

This equation has a critical point: $v = v_c$; $r = r_c = GM_{\odot} / (2v_c^2)$

Integration gives a family of curves:

$$\left(\frac{v}{v_c^2} \right)^2 - \ln \left(\frac{v}{v_c^2} \right)^2 = 4 \ln \frac{r}{r_c} + \frac{2GM_{\odot}}{rv_c^2} + C$$

Five regimes of solutions



Solution IV through the critical point ($\nabla \cdot \mathbf{C} = -3$) is the solar wind:
subsonic near the Sun \rightarrow supersonic beyond the critical point

The interplanetary magnetic field (IMF)

Close to the Sun:

Flow is radial, \mathbf{B} is nearly radial
Assume, that flow remains radial

\mathbf{B} is frozen-in to the rotation surface
and in the outflowing plasma

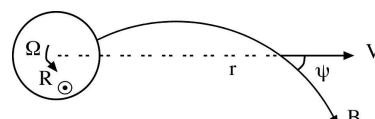
Flow speed $\wedge \mathbf{B}$: $V_\wedge = V \sin \Upsilon$

"Speed of the fl." $\wedge \mathbf{r}$: $W(r - R_\odot)$

$$V \sin \Psi = \Omega(r - R_\odot) \cos \Psi \Rightarrow \tan \Psi = \frac{\Omega(r - R_\odot)}{V}$$

For large r :

Archimedes' spiral

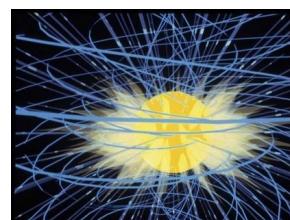


$\approx 44^\circ$ at Earth (1 AU)
 $\approx 57^\circ$ at Mars (1.5 AU)
 $\approx 88^\circ$ at Neptune (30 AU)

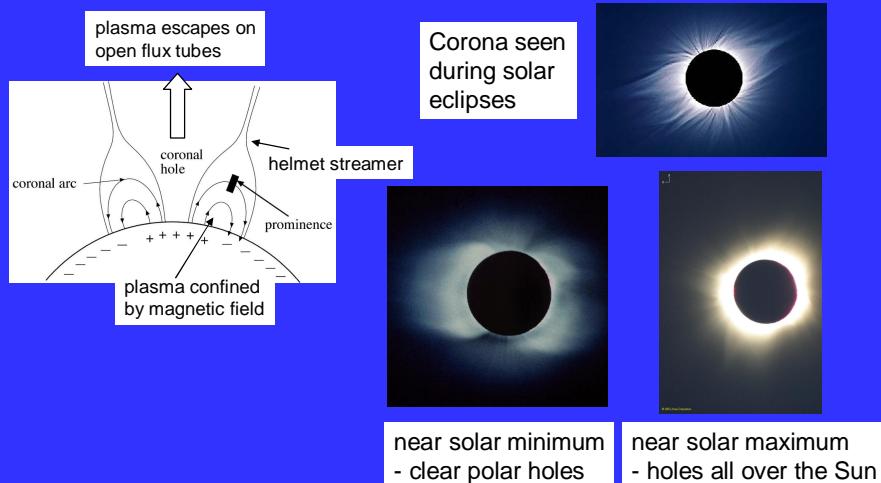
Far from the Sun:

$B \rightarrow r^{-1}$ In the ecliptic (tight spiral)

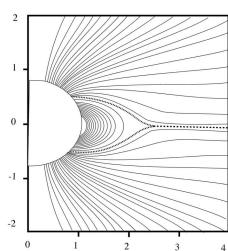
$B \rightarrow r^{-2}$ In the polar directions



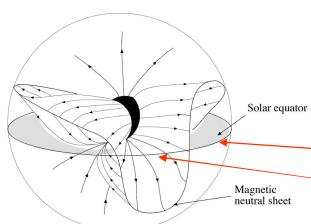
Sources of the solar wind: Coronal holes



How does the real solar wind look like?



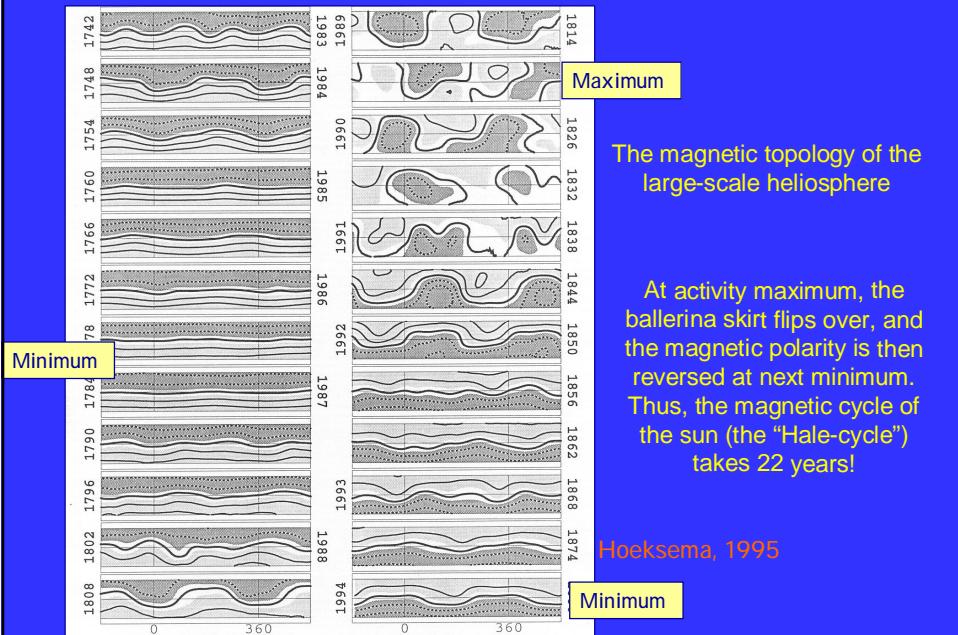
Formation of the
heliospheric current sheet
(Pneuman and Knopp)



The real current sheet
is curved
“Ballerina skirt”
(Alfvén)

The Earth can be
“toward” sector or in
“away” sector

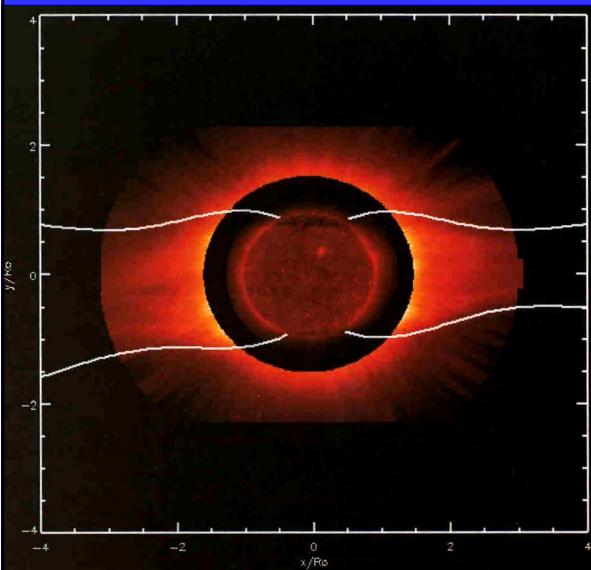
The ballerina dancing through the solar cycle



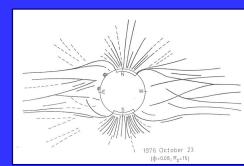
The magnetic topology of the large-scale heliosphere

At activity maximum, the ballerina skirt flips over, and the magnetic polarity is then reversed at next minimum. Thus, the magnetic cycle of the sun (the "Hale-cycle") takes 22 years!

The equatorial streamer belt



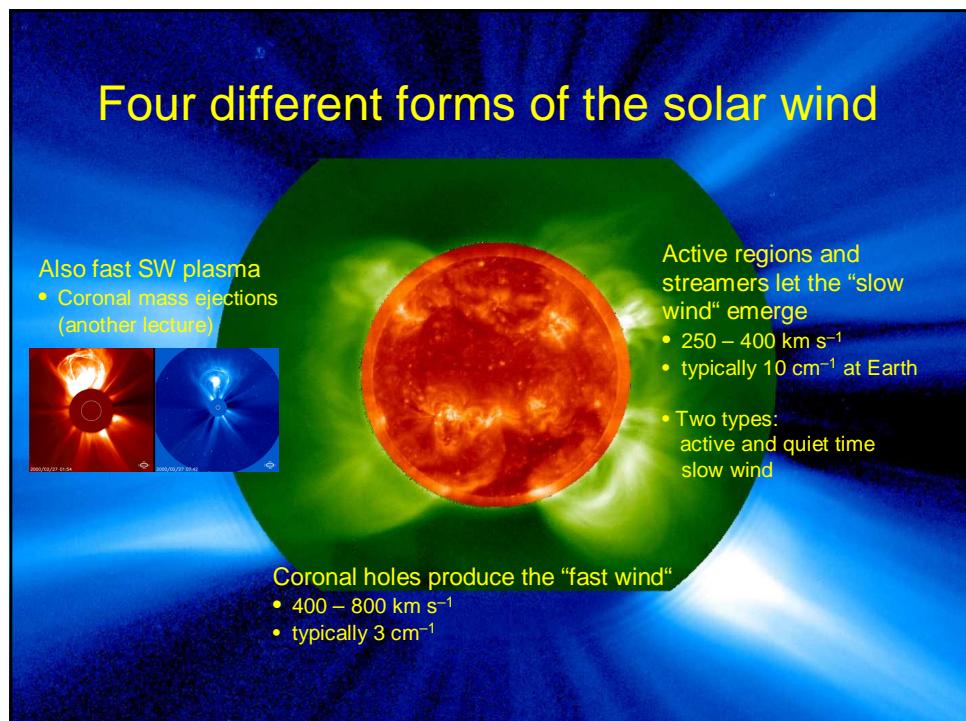
The boundaries of coronal holes and the streamer belt, as seen by EIT and UVCS on SOHO



Solar wind at 1 AU

	minimum	average	maximum
v	200 km/s	400 km/s	900 km/s
n	$4 \times 10^5 \text{ m}^{-3}$	$6.5 \times 10^6 \text{ m}^{-3}$	10^8 m^{-3}
T_e	$5 \times 10^3 \text{ K}$	$2 \times 10^5 \text{ K}$	10^6 K
T_p	$3 \times 10^3 \text{ K}$	$5 \times 10^4 \text{ K}$	10^6 K
B	0.2 nT	6 nT	80 nT
v_A	30 km/s	60 km/s	150 km/s

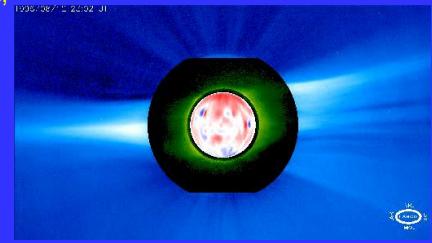
But average is a pretty useless concept here!



The two basic types of solar wind

1. Fast wind in high speed streams

High speed	$400\text{-}800 \text{ km s}^{-1}$	All values with reference to 1 AU
Low density	3 cm^{-3}	
Low particle flux	$2 \times 10^8 \text{ cm}^{-2} \text{ s}^{-1}$	
Helium content	3.6%, stationary	
Source	coronal holes	
Signatures	stationary for long times, all streams are alike, Alfvénic fluctuations	



2. Low speed wind of "interstream" type

Low speed	$250\text{-}400 \text{ km s}^{-1}$
High density	10.7 cm^{-3}
High particle flux	$3.7 \times 10^8 \text{ cm}^{-2} \text{ s}^{-1}$
Helium content	below 2%, highly variable
Source	helmet streamers near current sheet, at activity minimum (???)
Signatures	generally very variable, sector boundaries imbedded,

The two basic types of solar wind

Differences and similarities

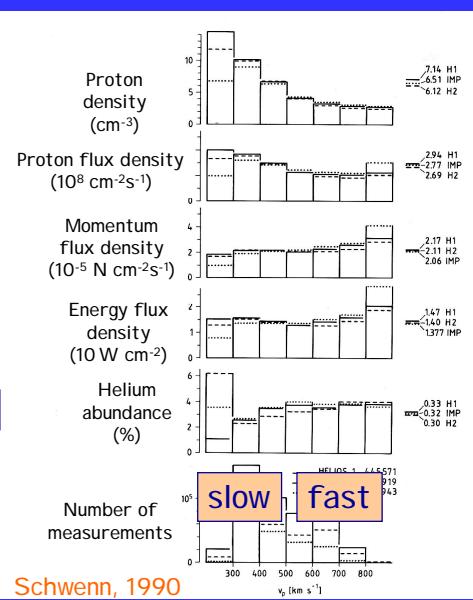
Strongly different

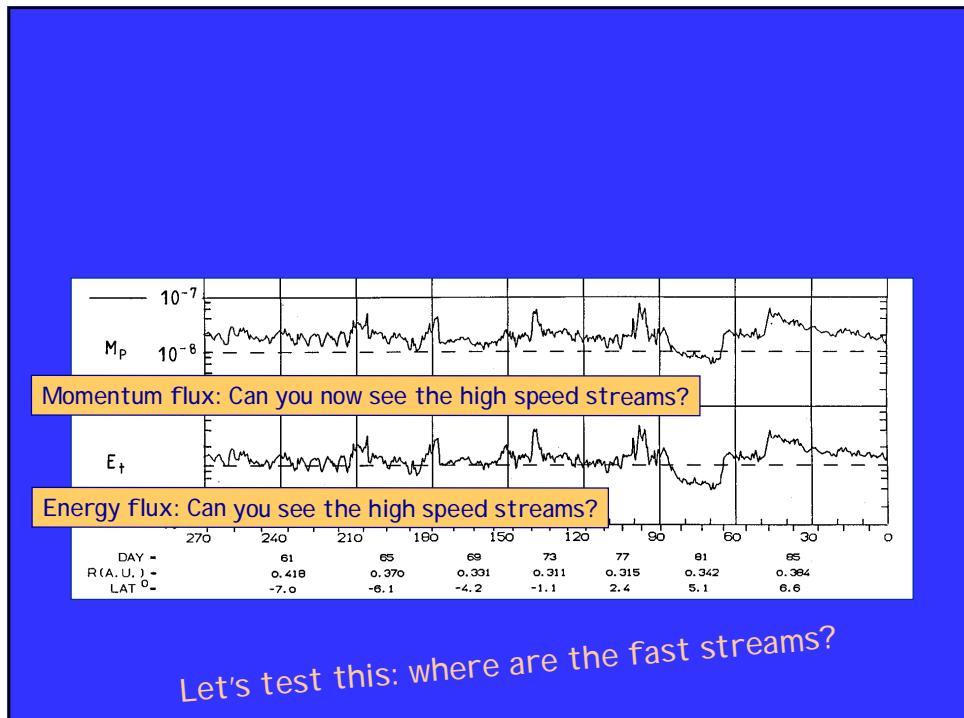
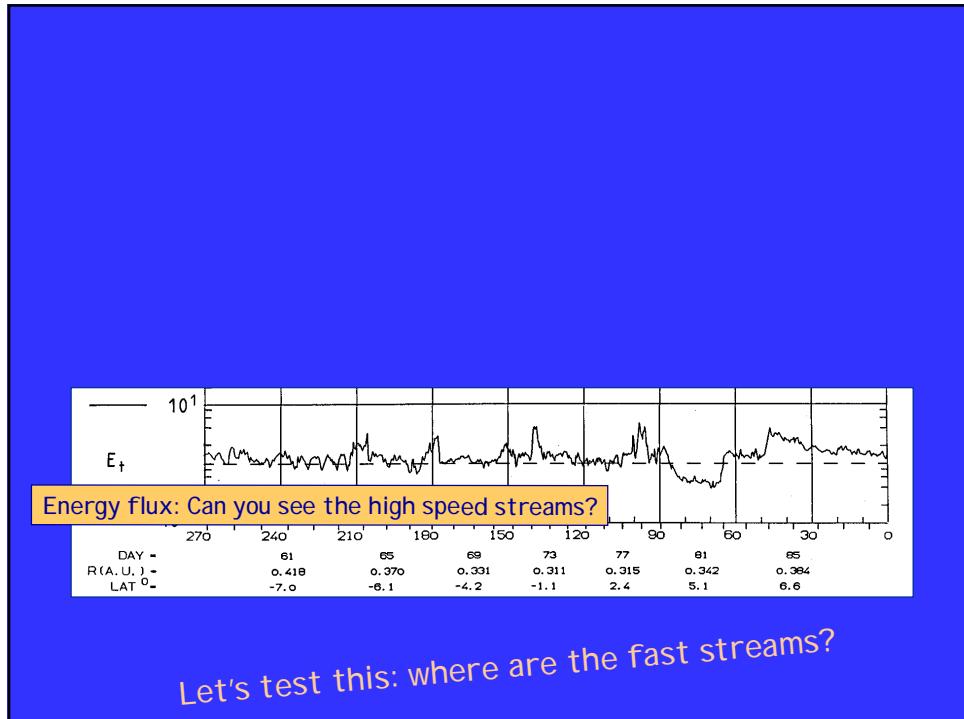
Significantly different!

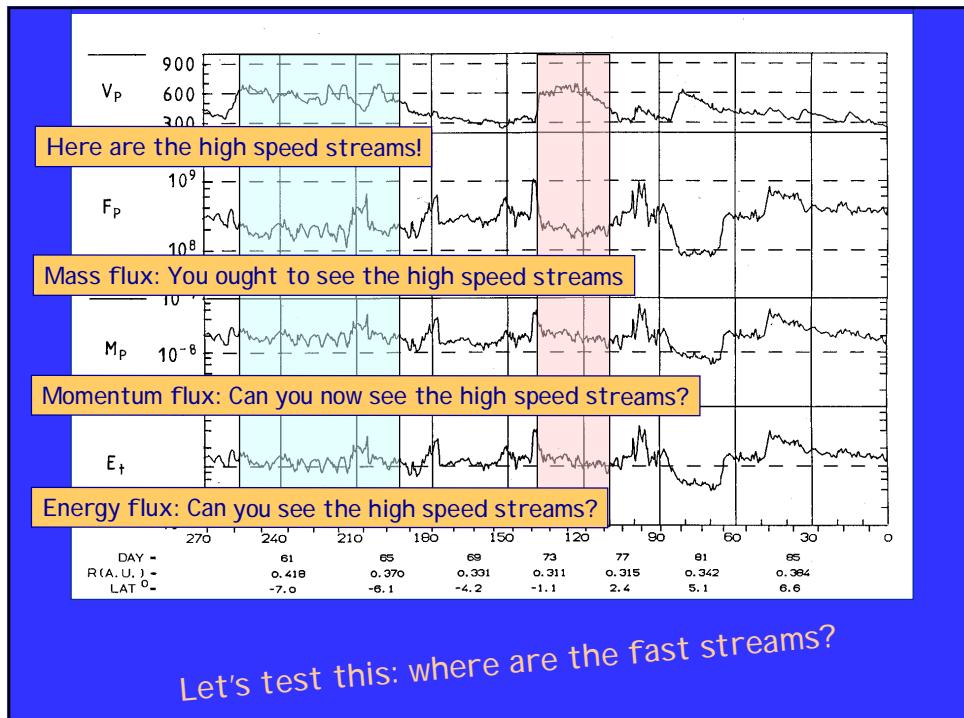
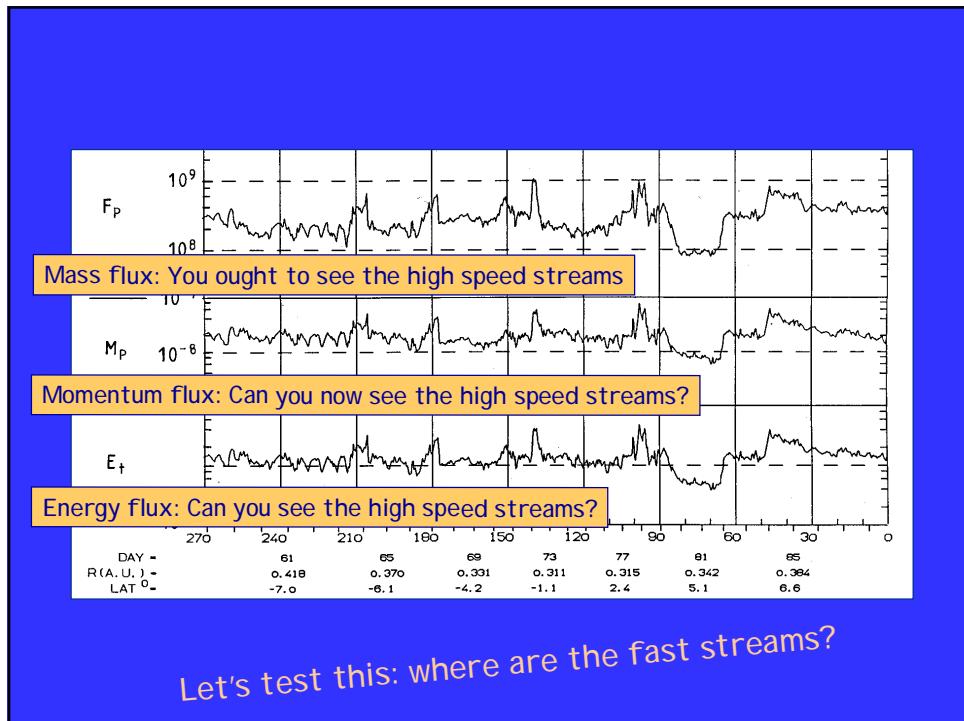
Almost invariant!

Also almost invariant!

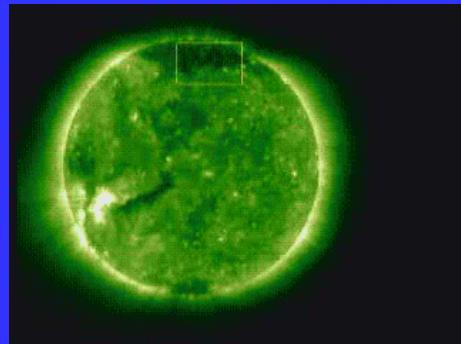
Significantly different!





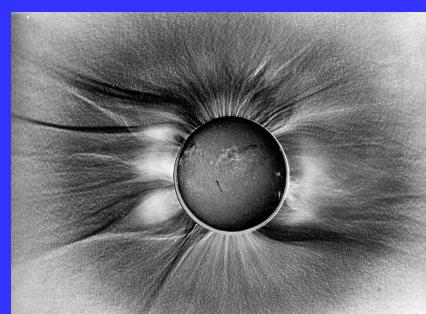
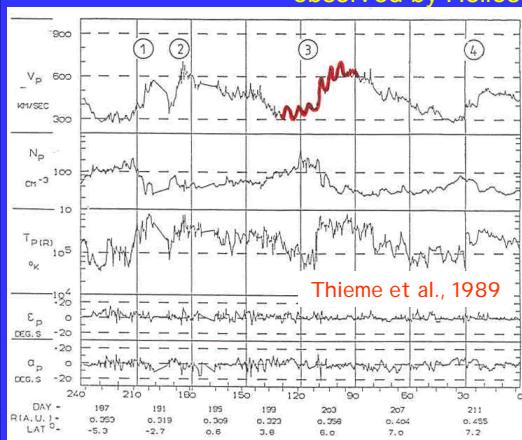


The “fast” solar wind, from high latitude coronal holes



The SOHO instruments have shown, that the fast solar wind from coronal holes emerges from the network boundaries, in particular from their intersections

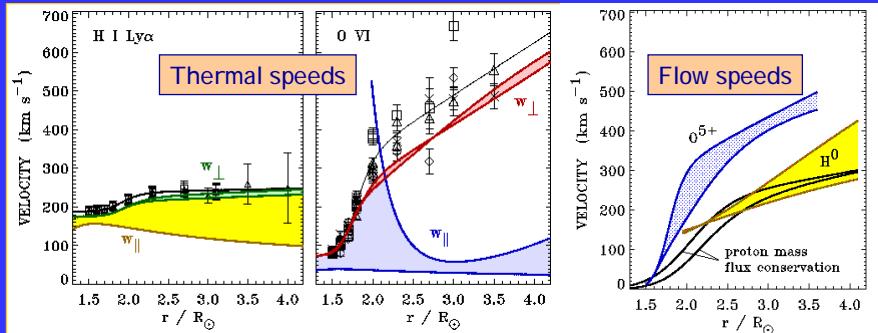
The filamentary structure of the solar wind observed by Helios near 0.3 AU



Koutchmy et al., 1978

The angular size of the solar wind flux tubes is $\sim 5^\circ$. It corresponds well to the scale size of the expanded polar plume structure known from eclipses.
At 1 AU these structures are usually flattened out.

The “fast” solar wind, from high latitude coronal holes

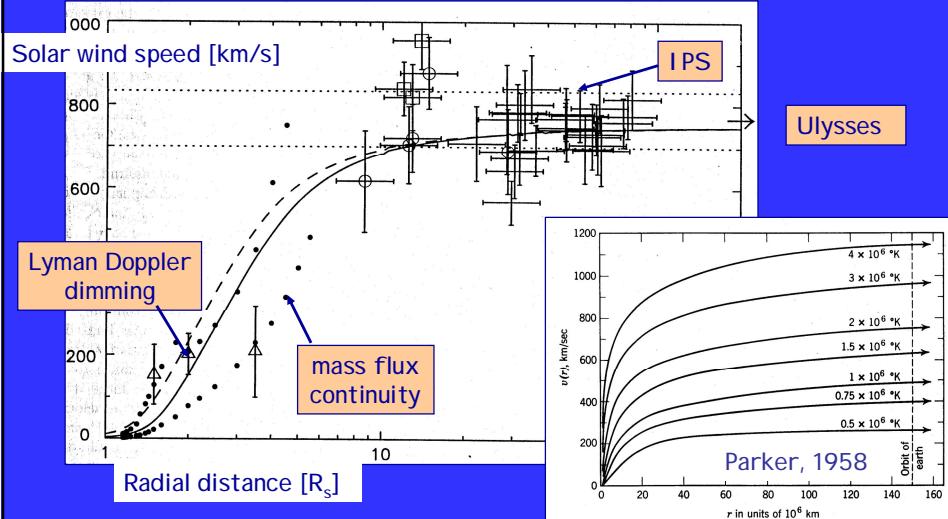


UVCS on SOHO measured the proton and O^{5+} ion thermal speeds in a coronal hole. For reference, the electron thermal speeds are shown as well.

Note the pronounced anisotropy and acceleration of oxygen with respect to the protons, thus indicating ion cyclotron heating.

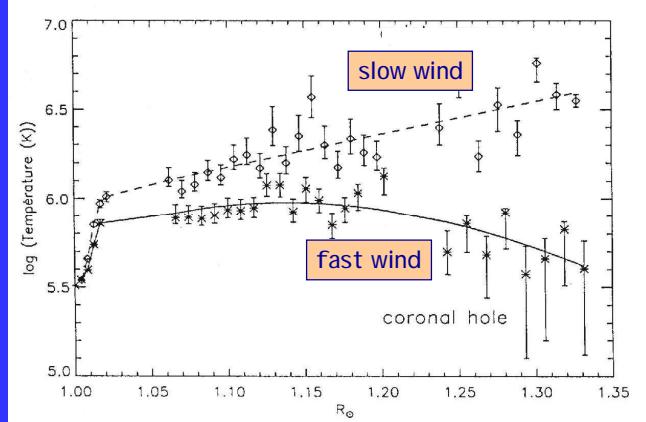
Note that the fast wind is almost “ready” by $3 R_{\text{sl}}$

Fast solar wind speed profile



Note: the fast solar wind has substantial speed as low down as $3 R_s$ and reaches its final speed inside $10 R_s$

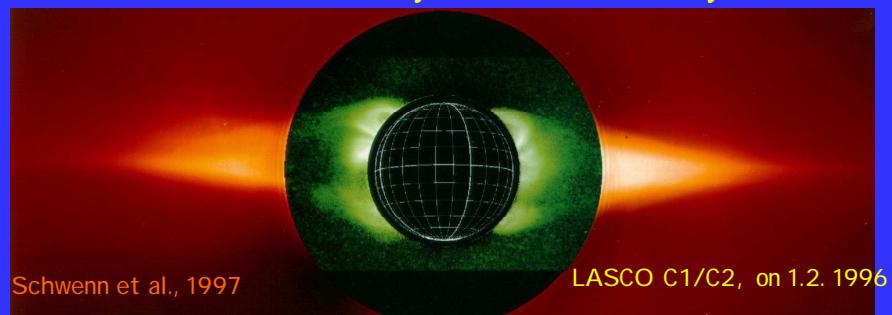
The fast solar wind, from high latitude coronal holes



CDS/SUMER: profiles of electron temperature from EUV line-ratios

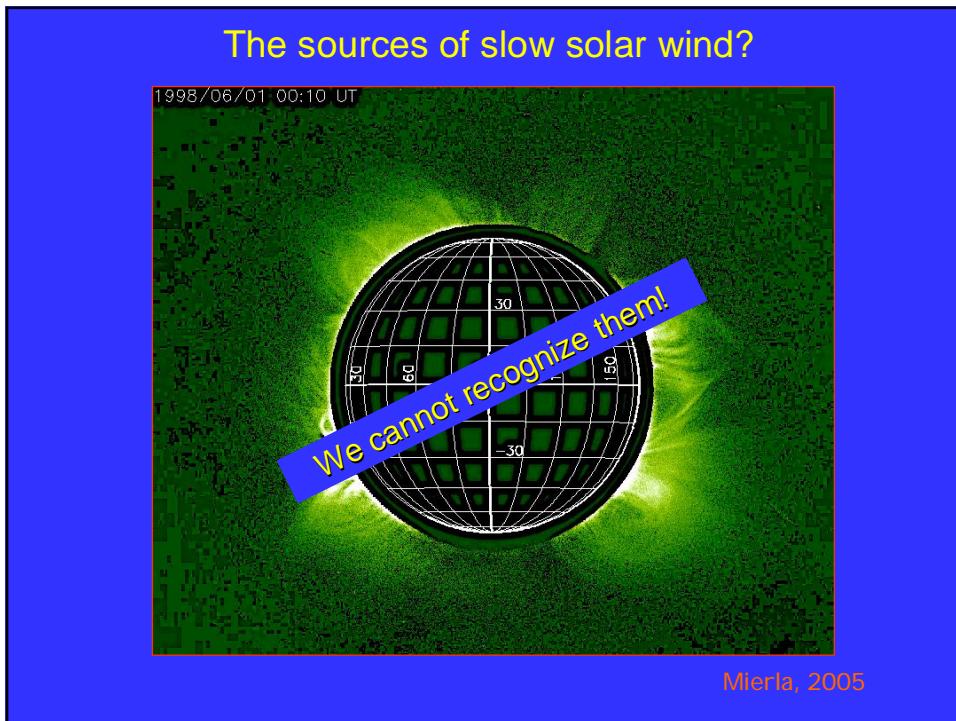
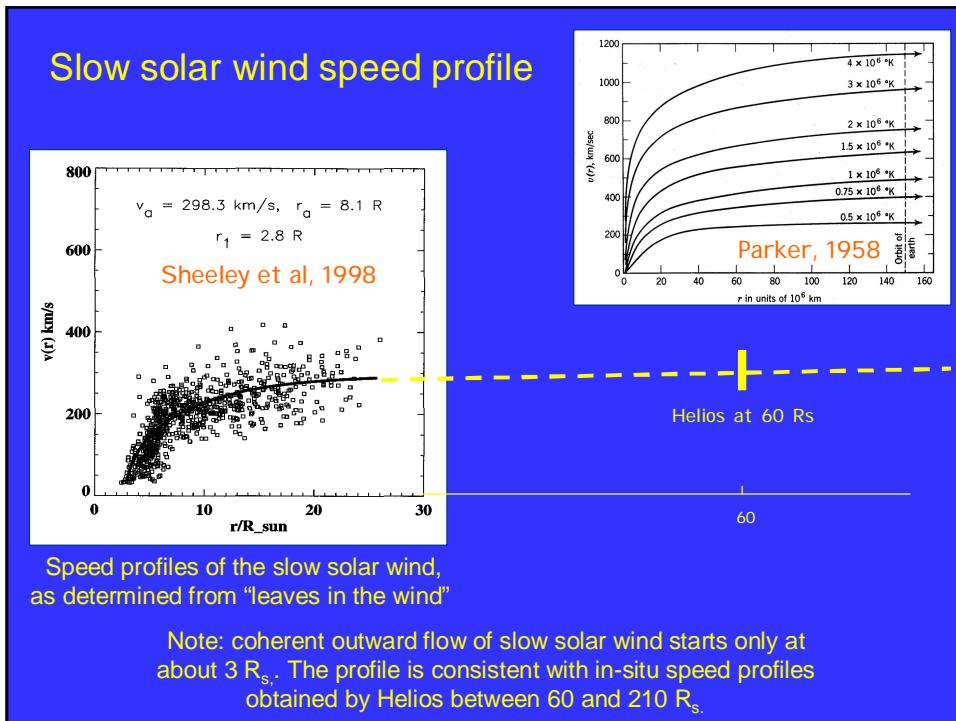
- The electron temperature is generally lower and drops off more rapidly in coronal holes than in streamers!
- In fast wind electrons are cooler than protons!

The corona at activity minimum in early 1996

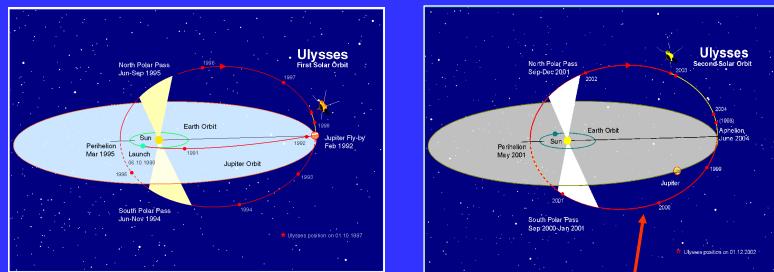


- There are magnetic multipole structures at mid-latitudes, in addition to the general dipole,
- these helmets may involve multiple current sheets,
- the mid latitude loops appear to be very stable in time, i.e., they extend over substantial longitudes as do the underlying photospheric neutral lines,
- the near-equatorial helmets vary strongly and are often absent.

The streamer sheet (only 30 deg wide in interplanetary space!)
and the HCS imbedded in it are products of the mid-latitude
streamers close to the Sun, NOT of the activity belt near the equator!

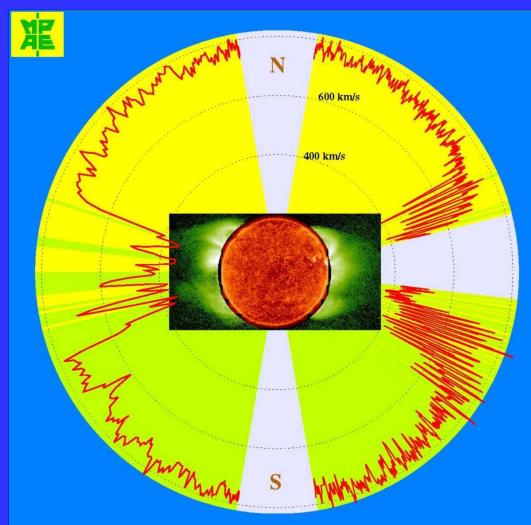


Ulysses observations



Present location January, 2006
The third and last revolution
around the Sun

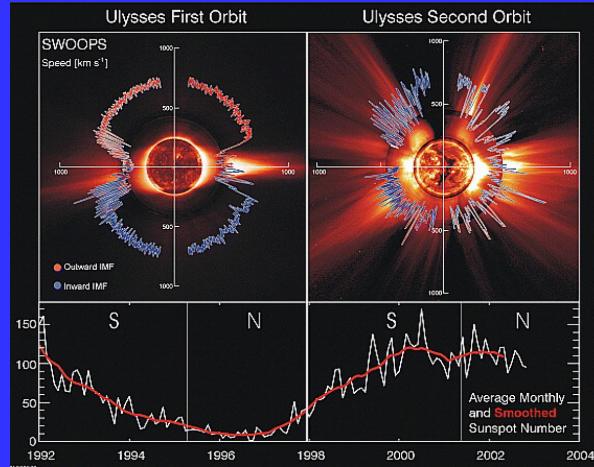
The 3D solar wind at activity minimum



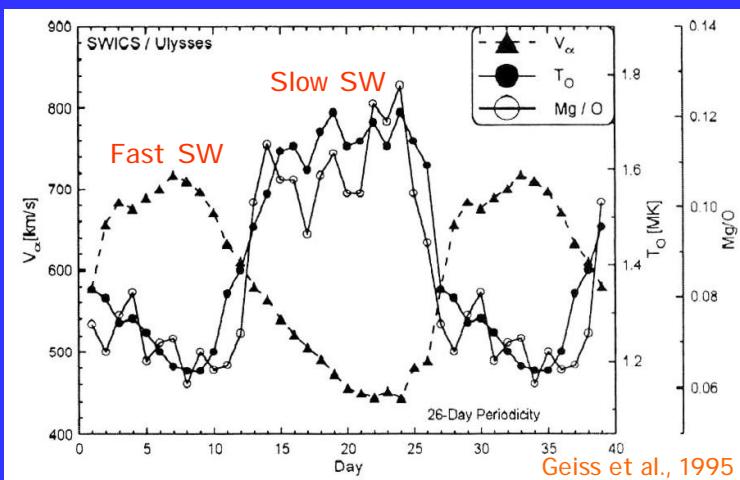
Ulysses was almost
permanently
encountering fast solar
wind, except from a
narrow, near-equatorial
belt of slow solar wind,
thus confirming earlier
measurements (e.g., from
IPS, Helios).

Ulysses observations of solar wind speed and
magnetic sector structure, observed during a full
orbit around solar activity minimum.

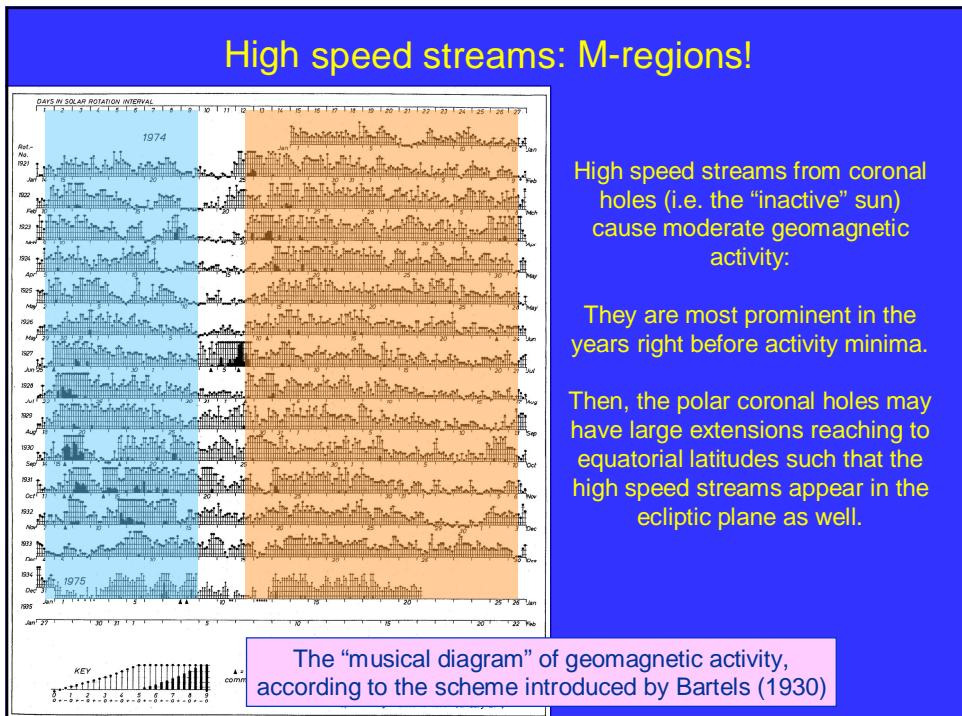
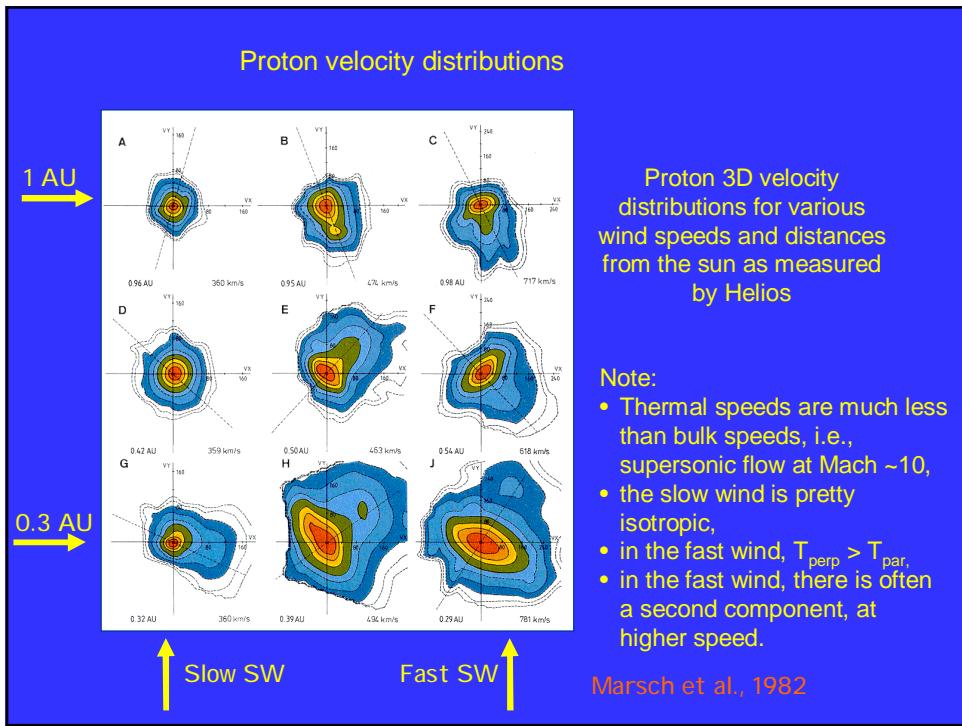
Minimum and maximum epochs
are very different!



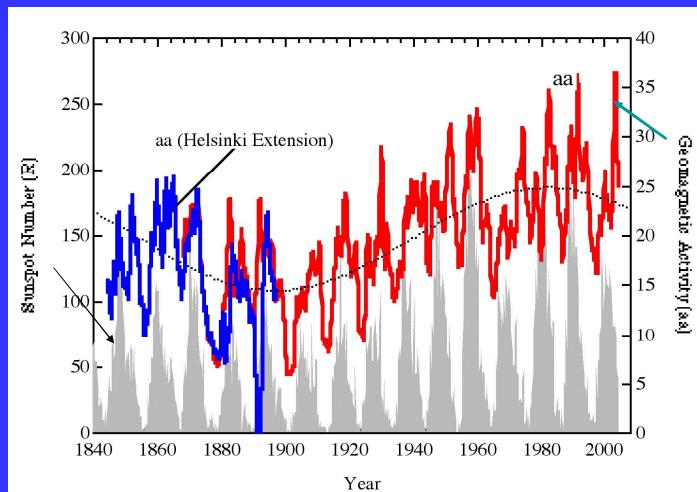
Solar wind element and charge state composition
and the large-scale stream structure



Signatures of stream structure in heavy ion
composition and charge state, measured from **Ulysses**.



Solar cycle and the Earth



Solar Wind Sources and their variations over the activity cycle

Key questions still waiting for answers:

- Two types of corona and solar wind: basic differences?
- Where and how is the slow solar wind emerging?
- How is the fast wind accelerated to its extreme speeds?
- Why that sharp boundaries between coronal and solar wind states?
- How to recognize CMEs and flares before they occur?
- 11 year activity and 22 year Hale cycles: significance for Sun-Earth system?

