



# Electromagnetic Fields Inside the Magnetosphere

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## Outline

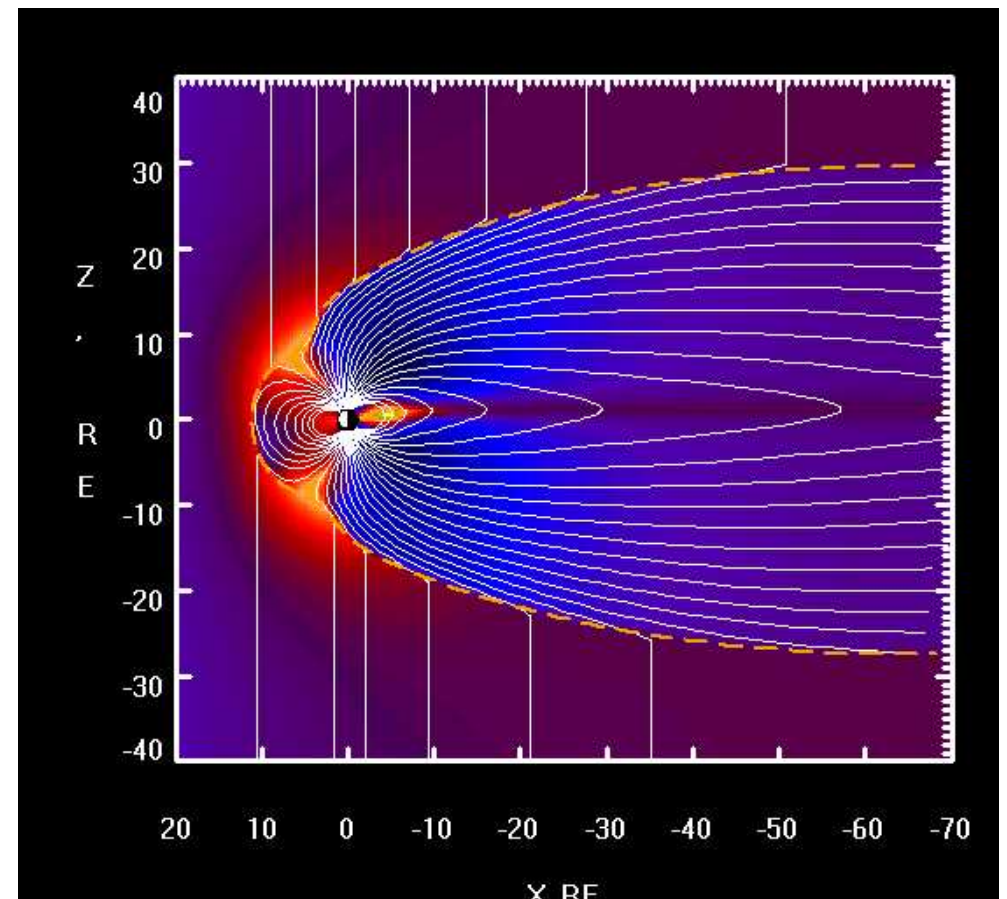
- Introduction to large-scale electromagnetic fields
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  - ▷ Magnetospheric convection
  - ▷ Magnetospheric current systems
  - ▷ Magnetospheric dynamics
- Motivation
- Magnetic field models
  - ▷ Historical models
  - ▷ Present-day models
- Electric field models
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  - ▷ Ionospheric models
- Modelling methods
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  - ▷  $\mathbf{E} \cdot \mathbf{B} = 0$
- Deformation method
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  - ▷ Euler potentials
  - ▷ Theory
  - ▷ Example & Model construction
  - ▷ 3D electric fields
- Modelling example
  - ▷ SSC, April 06, 2000
  - ▷ Dayside compression



# Intro to Large-Scale Electromagnetic Fields

## Magnetic Field Geometry

- Internal magnetic field
- Dipole tilt
- Solar wind pressure
- Open field lines
- Closed field lines
- Cusp



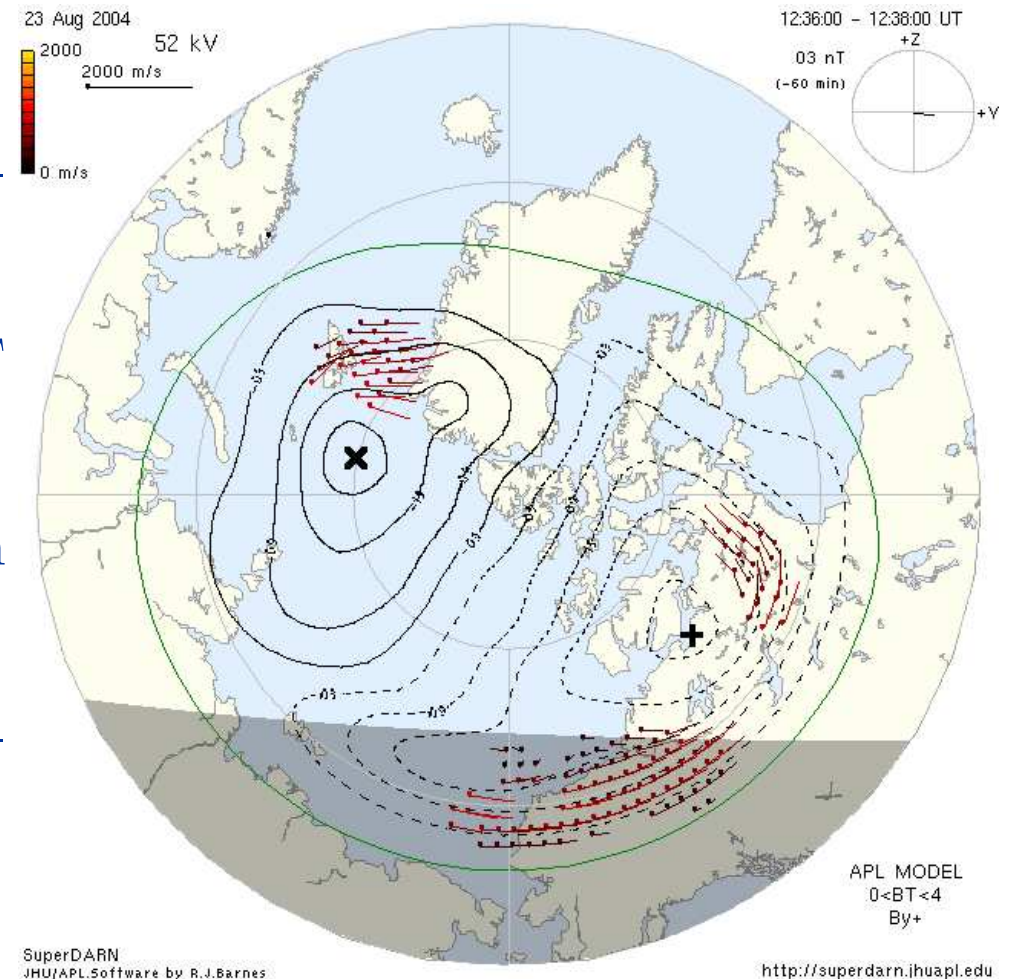
<http://modelweb.gsfc.nasa.gov/magnetos/data-based/modeling.html>



## Intro to Large-Scale Electromagnetic Fields (cont.)

### Convection

- Ionosphere:
- Global convection maps to the ionospheric convection
- Typically two-shell pattern (IMF  $B_z < 0$ )
- High-latitude convection from noon to midnight
- Return convection at lower latitudes via dawn and dusk
- Can be measured by radars



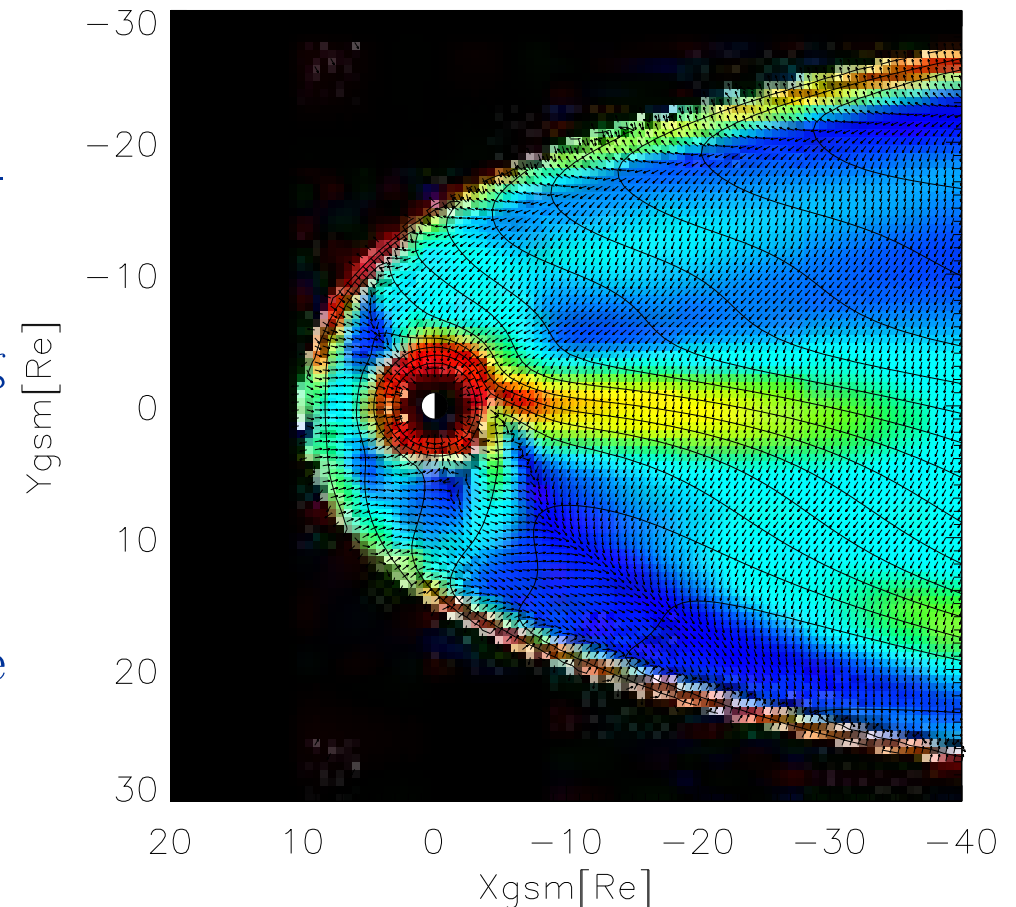
<http://superdarn.jhuapl.edu/>



## Intro to Large-Scale Electromagnetic Fields (cont.)

### Convection (cont.)

- Equatorial plane:
- Earthward convection in the mid-tail
- Plasma flow around the co-rotating plasmasphere to the dayside
- Return convection along the flanks
- Electric field associated with the convection  $\sim 0.5\text{mV/m}$
- Cannot be globally measured

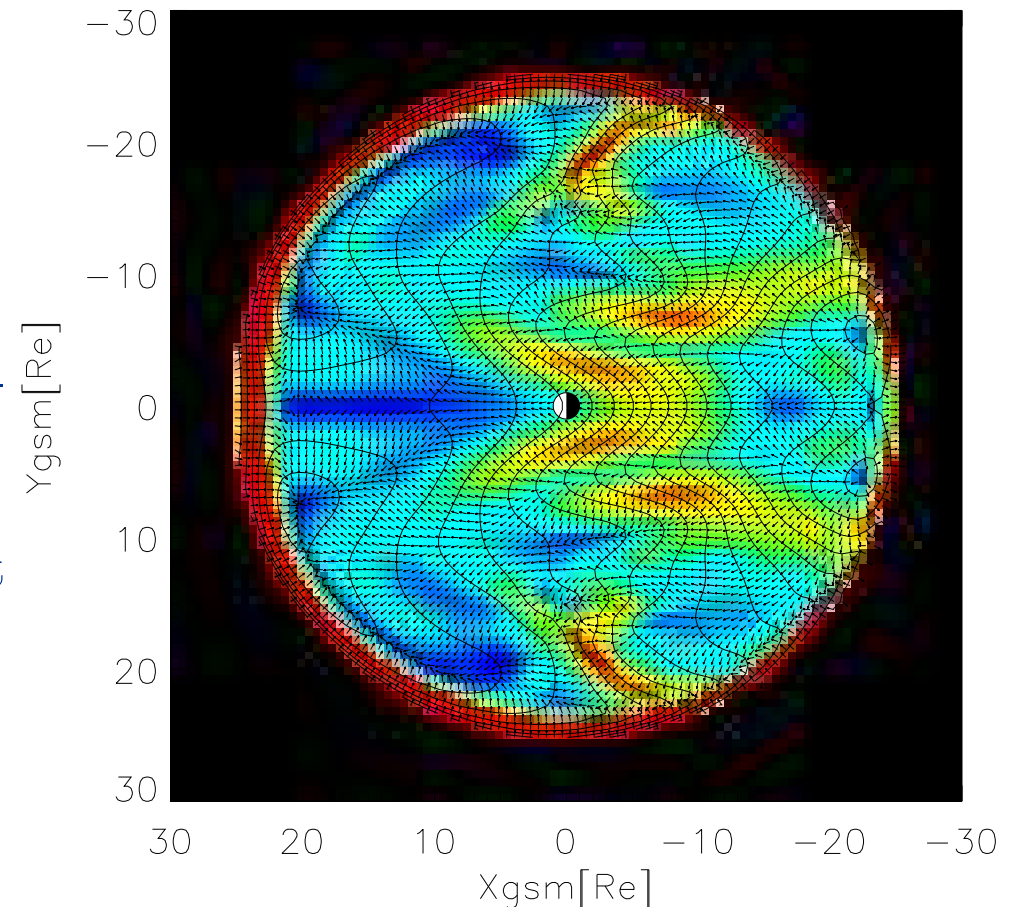




## Intro to Large-Scale Electromagnetic Fields (cont.)

### Convection (cont.)

- Tail cross-section:
- Convection across the tail lobes towards the plasma sheet
- Turns to earthward convection at the plasma sheet



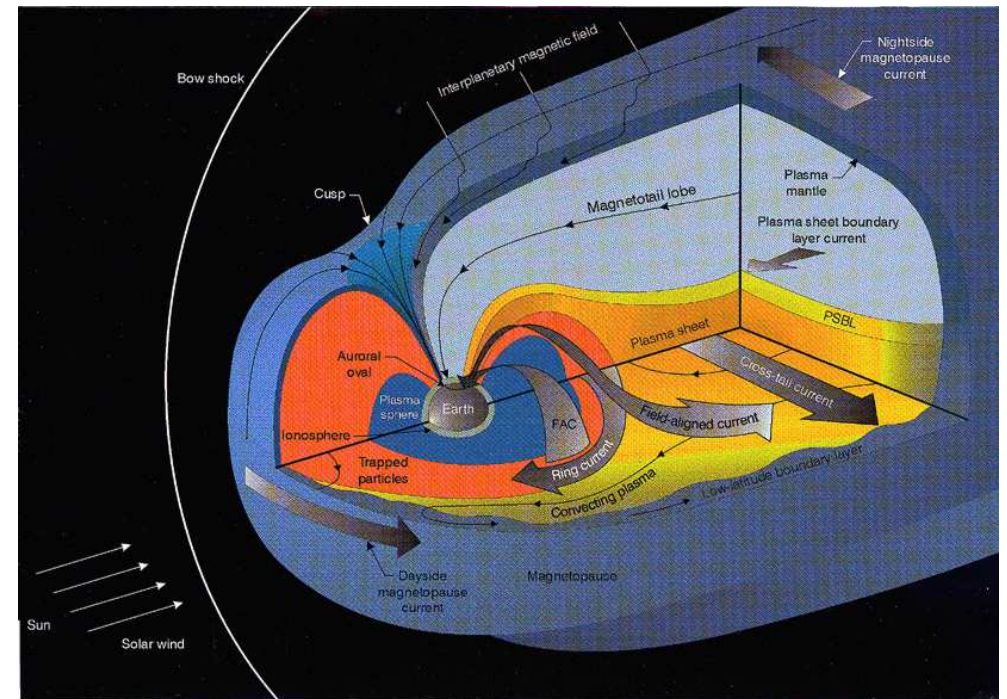




## Intro to Large-Scale Electromagnetic Fields (cont.)

### Current Systems

- Cross-tail current
- Ring current
- Magnetopause current
- Field-aligned currents

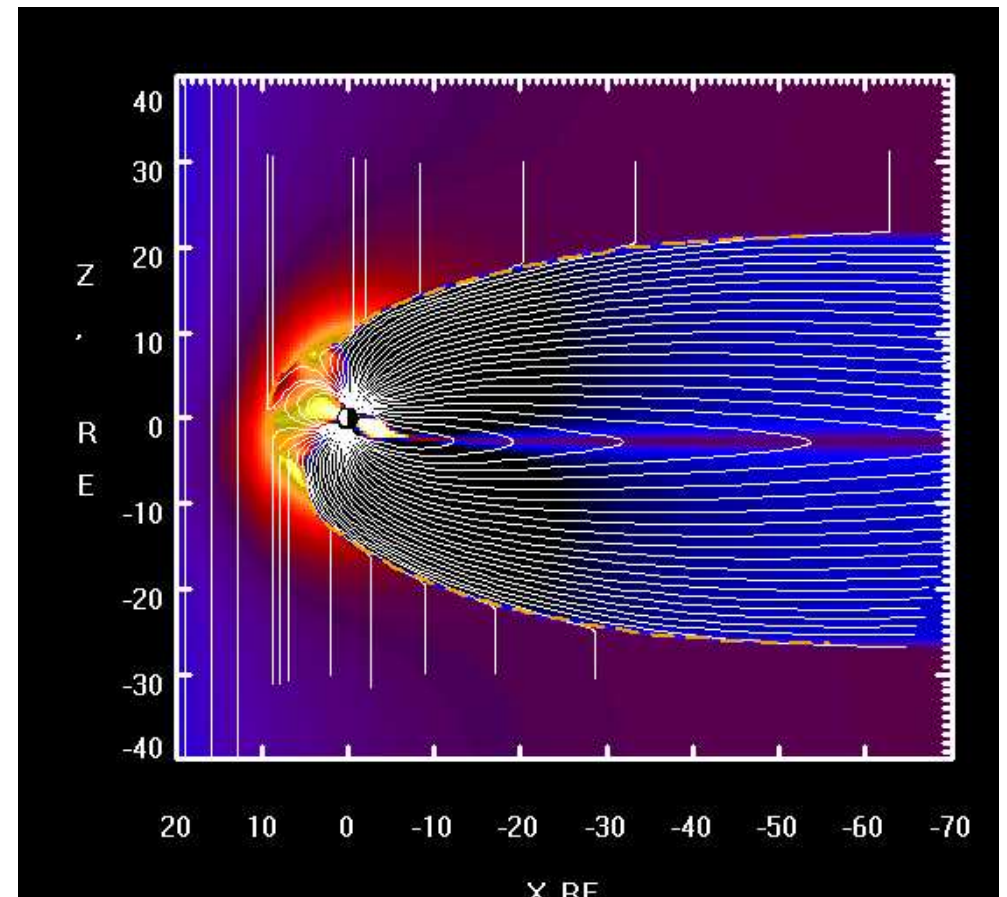




## Intro to Large-Scale Electromagnetic Fields (cont.)

### Magnetospheric Dynamics

- Magnetic storms
- Magnetopause compression
- Ring current enhancement
- Substorms
- Tail stretching
- Dipolarization



<http://modelweb.gsfc.nasa.gov/magnetos/data-based/modeling.html>



## Why Analytical Field Models?

### Physical Reasons

- Link between the interplanetary medium and the ionosphere
- Guide and Energize charged particles
- Confines radiation belts and controls the auroral phenomena
- Store a huge amount of energy

### Practical Reasons

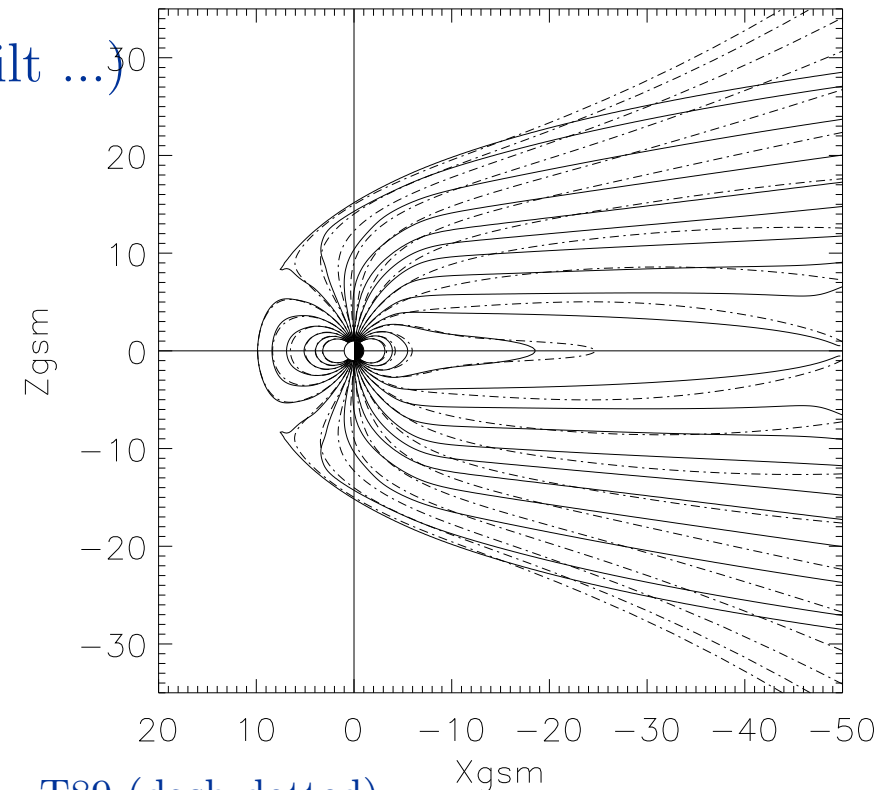
- Satellite conjunctions (along magnetic field lines)
- Ground-based stations and spacecraft conjunctions
- Test particle simulations (Dr. Ganushkina)
- Mapping of ionospheric convection to the magnetosphere
- Subtract the background field from observation





## Magnetic Field Models

- For historical models: See Stern, JGR, 1994.
- IGRF (International Geomagnetic Reference Field)
- Geopack (coordinate transformations, tilt ...)
- T89 (Tsyganenko, JGR, 1989)
  - ▷ Computationally fast
  - ▷ Extremely widely used
  - ▷ Parametrized by Kp
- T96
  - ▷ More realistic current systems
  - ▷ Larger set of input parameters



T89 (dash-dotted)

Toivanen (solid)

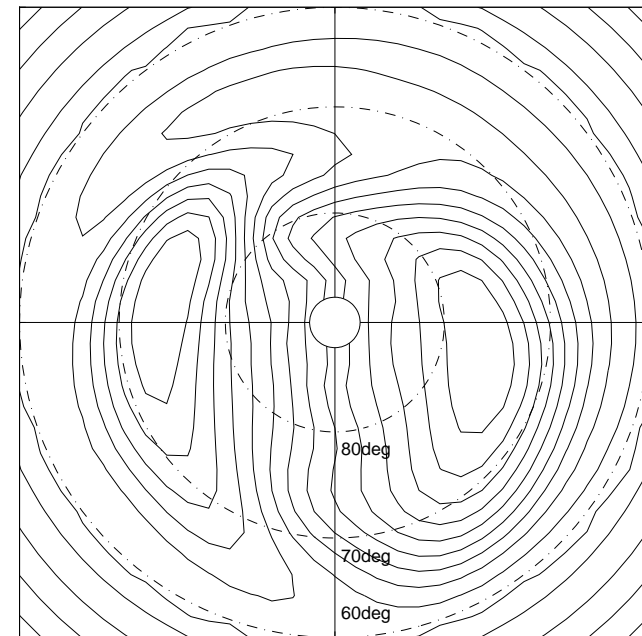
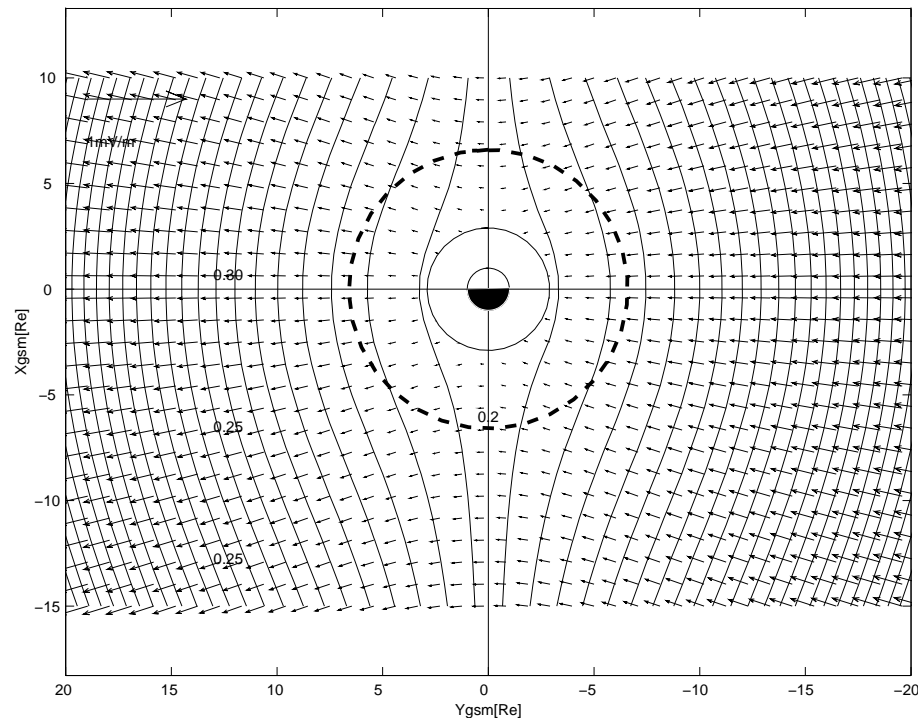
- All available at

<http://modelweb.gsfc.nasa.gov/magnetos/data-based/modeling.html>



## Electric Field Models

- Equatorial model: Volland, JGR, 1973.
  - ▷  $\varphi = Ar^\gamma \sin \phi$ , where  $A$  is a constant, and  $\gamma = 2$ , typically.
- Ionospheric model: Heppner and Maynard, JGR, 1987.





## Modelling Methods and Shortcomings

### Magnetic Field

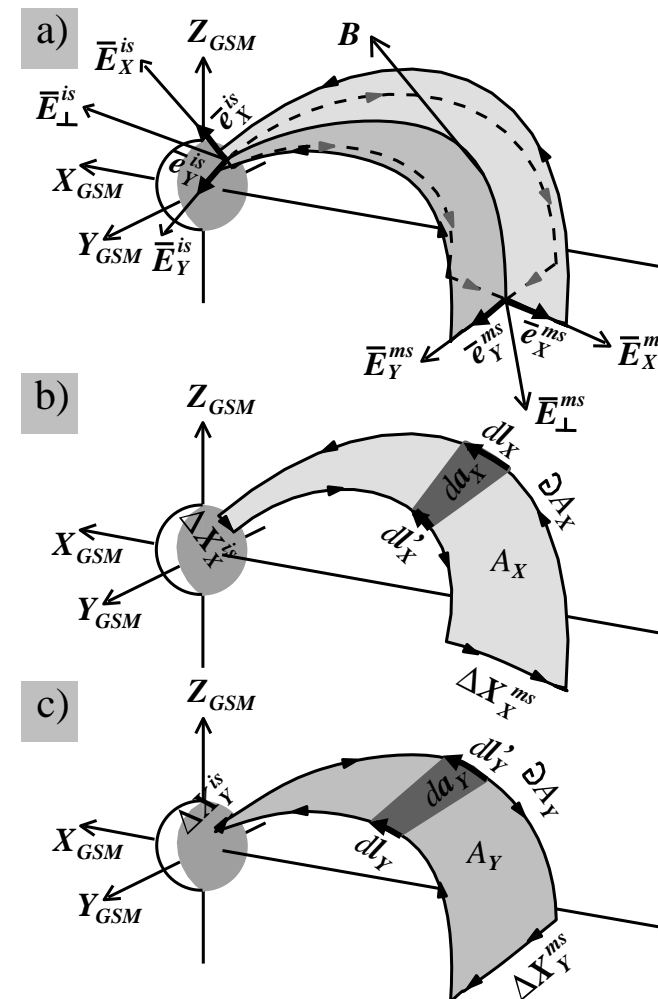
- Represent mathematically the field from each major current systems
- Represent the expected response of the field to physical factors that can be determined
- Calibrate the model against a database of average magnetic field observations, tagged by values of the solar wind parameters
- Force balance is not automatically built into the models
- Unreliable in highly time-variable situations such as substorms → event-oriented models (Deformation method below + Dr. Ganushkina)



## Modelling Methods and Shortcomings (cont.)

### 3D Electric Field

- Equipotential mapping (static)
- $\int_{\partial A} \mathbf{E} \cdot d\mathbf{l} = 0$
- Generalized mapping
- $\int_{\partial A} \mathbf{E} \cdot d\mathbf{l} = - \int_A \partial_t \mathbf{B} \cdot d\mathbf{a}$
- Induction included
- Computationally tedious
- Toivanen, P. K., et al., JGR, 1998.





## Deformation Method

### Motivation

- Solving Faraday's law is complicated ( $\nabla \times \mathbf{E} = -\partial_t \mathbf{B}$ )
- No explicit vector potential in magnetic field models
- Electrostatic component difficult to include
- $\partial_t \mathbf{A}$  may have a component parallel to  $\mathbf{B}$
- Euler potentials  $\rightarrow$  Physical electromagnetic fields

Field	$(\mathbf{A}, \phi)$	$(\alpha, \beta, \phi)$
$\mathbf{B}$	$\nabla \times \mathbf{A}$	$\nabla \alpha \times \nabla \beta$
$\mathbf{E}_{vacuum}$	$-\partial_t \mathbf{A} - \nabla \phi$	$-\partial_t(\alpha \nabla \beta) - \nabla \phi$
$\mathbf{E}_{plasma}$	-	$-\partial_t \alpha \nabla \beta + \partial_t \beta \nabla \alpha - \nabla \phi$





## Deformation Method (cont.)

□ For a given  $\alpha(x, y)$



□  $\hat{\alpha}(x, y) = \alpha(\hat{x}, \hat{y})$



- As to Homer, this can be done to a magnetic field:
- $\mathbf{B}' = \nabla \hat{\alpha} \times \nabla \hat{\beta} \Rightarrow \mathbf{B}' = \mathbf{T} \hat{\mathbf{B}}$
- Take a magnetic field given in some coordinates.
- Replace these coordinates by some transformed coordinates.
- Multiply the resulted magnetic field by the matrix  $\mathbf{T}$  (given on the next slide).
- The given field is then given in terms of the original coordinates
- References
  - ▷ Stern, D. P., JGR, 1987.
  - ▷ Tsyganenko, N. A., JGR, 1998.



## Generalized Deformation Method

□ Consider a transformation

$$\hat{q}_i = f_i(q_1, q_2, q_3)$$

□ Let  $\mathbf{A}' = \hat{\alpha} \nabla \hat{\beta}$

$$\Rightarrow \mathbf{A}' = \mathbf{M} \hat{\mathbf{A}}; M_{hl} = \frac{\hat{h}_l}{h_h} \partial_h f_l$$

□ Let  $[\nabla \phi]' = \nabla \hat{\phi}$

$$\Rightarrow [\nabla \phi]' = \mathbf{M} [\nabla \phi]$$

□ Let  $\mathbf{B}' = \nabla \hat{\alpha} \times \nabla \hat{\beta}$

$$\Rightarrow \mathbf{B}' = \mathbf{T} \hat{\mathbf{B}}; T_{ij} = \frac{1}{2} \varepsilon_{ihk} M_{hl} M_{kq} \varepsilon_{lqj}$$

□ Let  $\mathbf{E}' = -\partial_t \hat{\alpha} \nabla \hat{\beta} + \partial_t \hat{\beta} \nabla \hat{\alpha} - \nabla \hat{\phi}$

$$\Rightarrow \mathbf{E}' = \mathbf{M} (\hat{\mathbf{E}} + \partial_t \hat{\mathbf{q}} \times \hat{\mathbf{B}})$$

□ For deformed fields:

$$\square \mathbf{A}' \cdot \mathbf{B}' = 0$$

$$\square \nabla \cdot \mathbf{B}' = 0$$

$$\square \nabla \times \mathbf{B}' \neq 0, \text{ even if } \nabla \times \mathbf{B} = 0$$

$$\square [\nabla \phi]' \cdot \mathbf{B}' = 0, \text{ if } \nabla \phi \cdot \mathbf{B} = 0$$

$$\square \mathbf{E}' \cdot \mathbf{B}' = 0$$

$$\square \nabla \times [\nabla \phi]' = 0$$

$$\square \nabla \times \mathbf{E}' = -\partial_t \mathbf{B}'$$



## Initial Dipole Fields

□ In spherical coordinates  $(r, \theta, \varphi)$ :

□ Euler potentials:  $\alpha \propto r^{-1} \sin^2 \theta$  and  $\beta \propto \varphi$

$$\Rightarrow \mathbf{B} = \frac{\mu_0 m}{2\pi r^3} \cos \theta \mathbf{e}_r + \frac{\mu_0 m}{4\pi r^3} \sin \theta \mathbf{e}_\theta$$

$$\Rightarrow \mathbf{A} = \frac{\mu_0 m}{4\pi r^2} \sin \theta \mathbf{e}_\varphi$$

□ Assume equipotential dipole field lines:  $-\nabla \phi \cdot \mathbf{B} = 0$

$$\Rightarrow \phi(r, \theta, \varphi) = \phi_{is}(r_{is}, \theta_{is}(r, \theta), \varphi_{is}(\varphi))$$

□ Dipole magnetic field lines

$$\Rightarrow (r_{is}, \sin \theta_{is}, \varphi_{is}) = (\text{constant}, \left[ \frac{r_{is}}{r} \right]^{\frac{1}{2}} \sin \theta, \varphi)$$

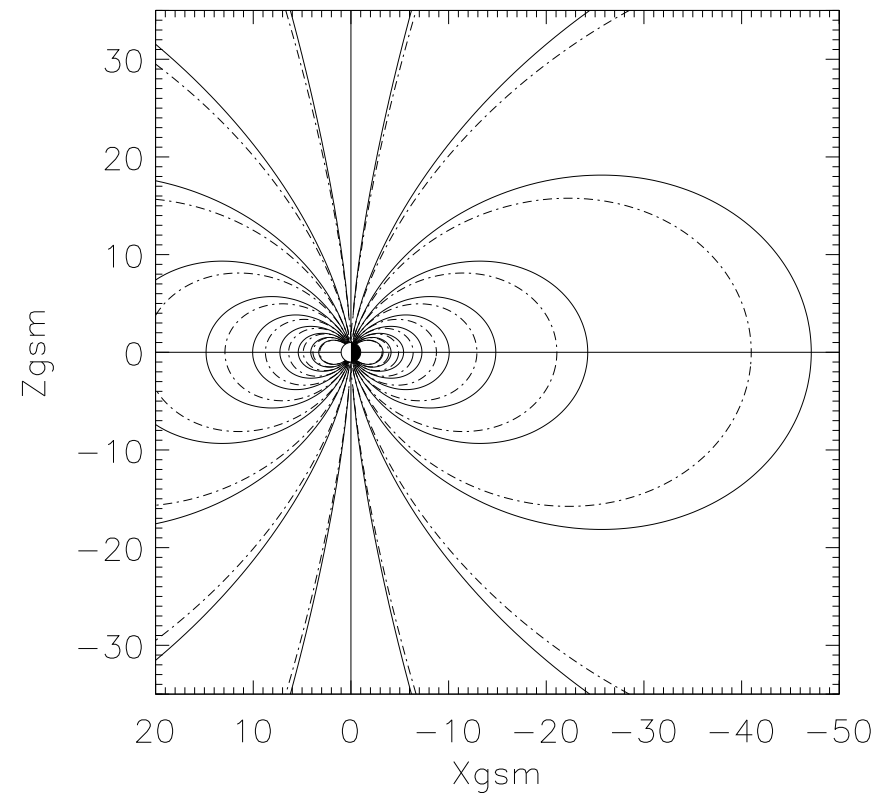
□ Assume an ionospheric electric field  $\mathbf{E}^{is} = (E_\theta^{is}, E_\varphi^{is})$

$$\Rightarrow \mathbf{E} = -\nabla \phi = -\nabla \phi_{is} = \left( \frac{r_{is}}{r} \right)^{\frac{3}{2}} \left( -\frac{1}{2} \frac{\sin \theta}{\cos \theta_{is}} E_\theta^{is} \mathbf{e}_r + \frac{\cos \theta}{\cos \theta_{is}} E_\theta^{is} \mathbf{e}_\theta + E_\varphi^{is} \mathbf{e}_\varphi \right)$$



## An Example: Dipole Expansion

- Dipole field (dash-dotted)
- $\alpha = \frac{\mu_o m_o}{4\pi r} \sin^2 \theta$
- $m \rightarrow m(r)$
- $\hat{r} = rH(r)^{-1}$
- $H(r) = a \left( 1 - \tanh \left( \frac{r-r_o}{\Delta r} \right) \right) + c$
- $a = (1 - c) \left( 1 - \tanh \left( \frac{r_{is}-r_o}{\Delta r} \right) \right)^{-1}$
- $m = m_o$ , if  $r = r_{is}$
- $m = cm_o$ , if  $r \rightarrow \infty$
- Expanded dipole (solid)
- $\Rightarrow$  Ring current

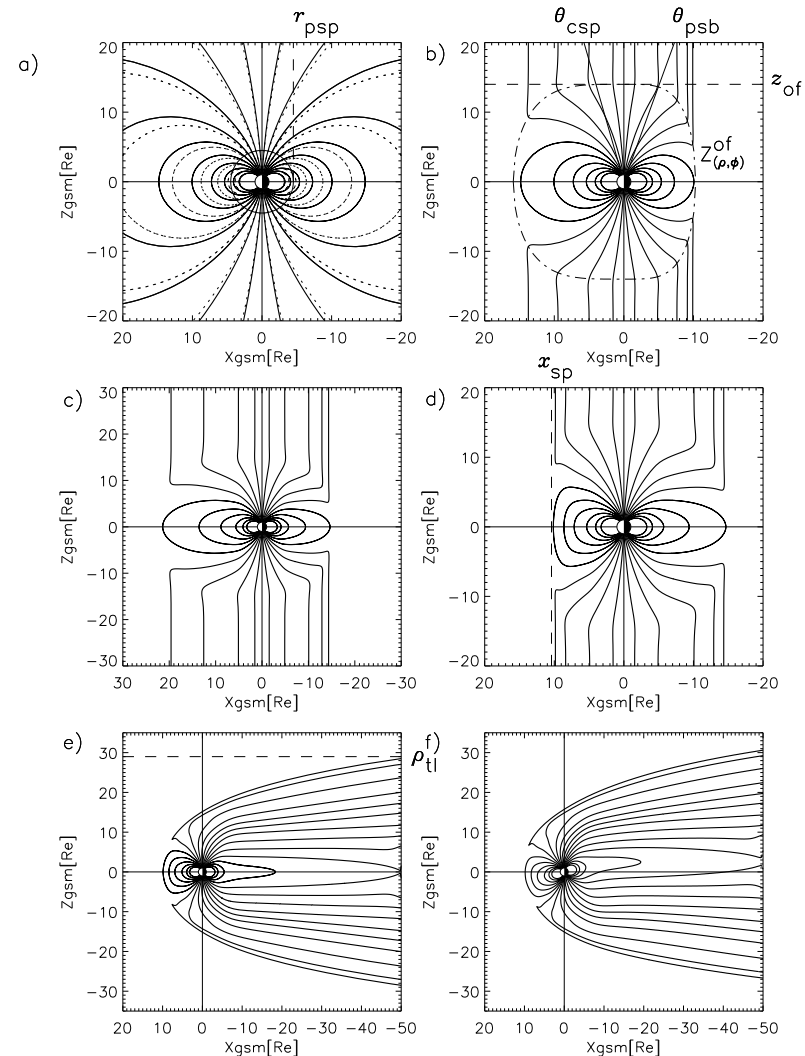




## Model Construction

### Noon-midnight meridian

- Consecutive deformations:
- Dipole expansion (a)
- Field topology (b)
- Near-Earth plasma sheet (c)
- Dayside compression (d)
- Magnetopause and lobes (e)
- Tilt effect (f)



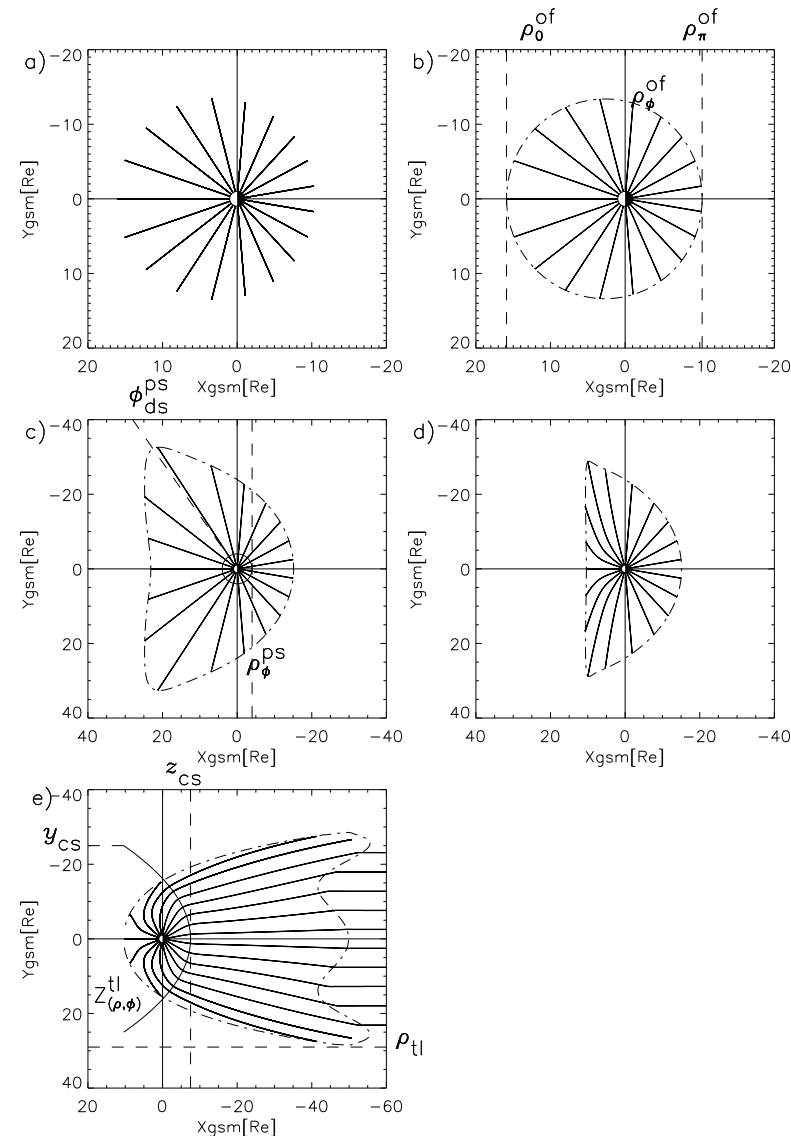




## Model Construction (cont.)

### Equatorial plane

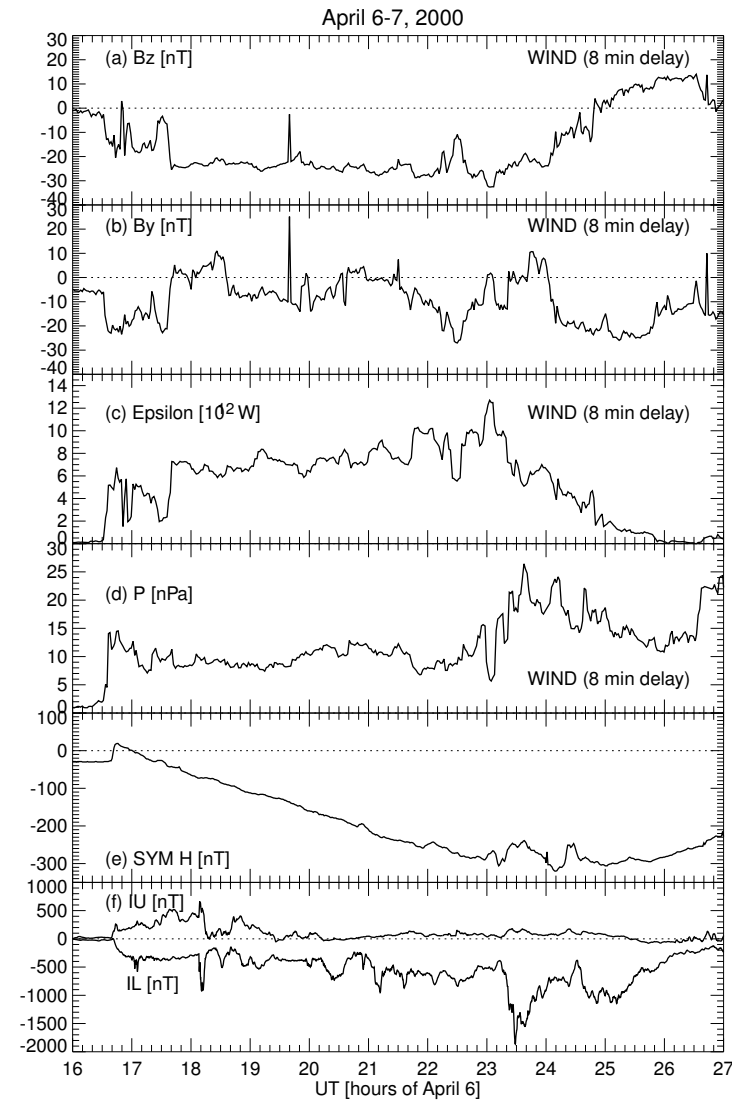
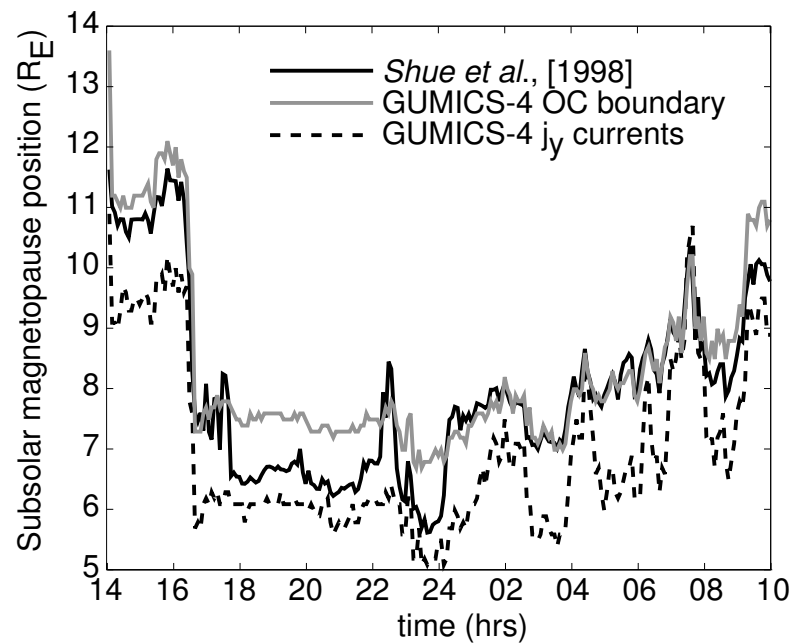
- Last closed field lines:
- Dipole expansion (a)
- Field topology (b)
- Near-Earth plasma sheet (c)
- Dayside compression (d)
- Magnetopause and lobes (e)





## SSC, April 06, 2000

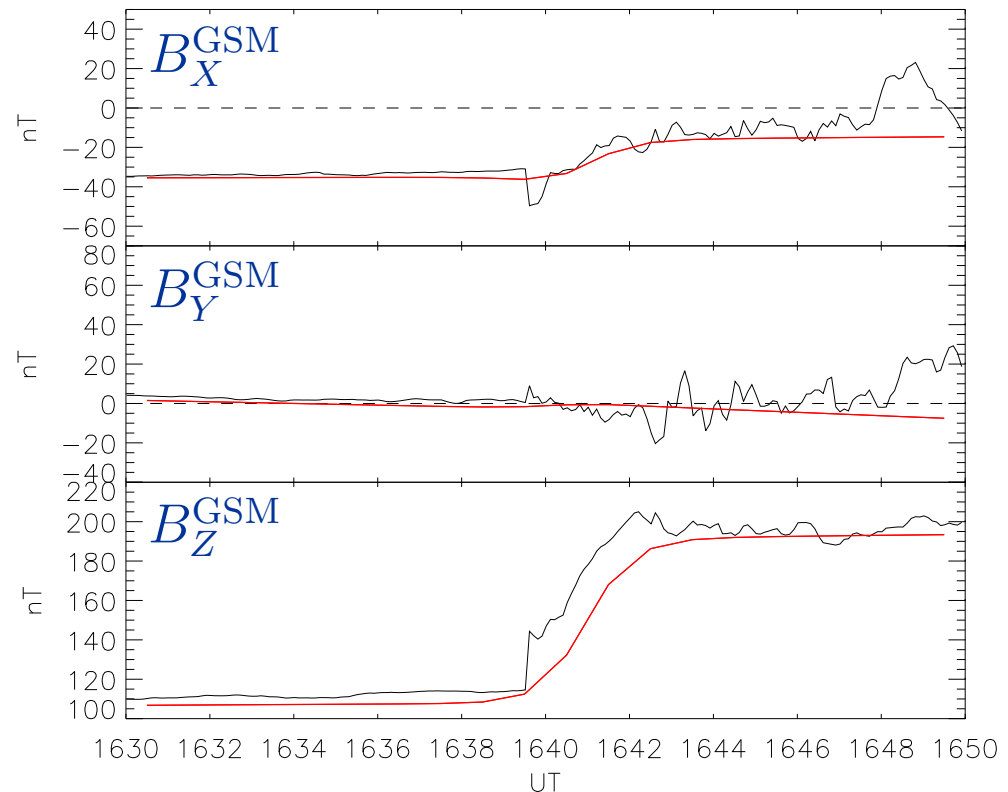
- Solar wind pressure pulse
- Negative excursion of IMF  $B_z$
- Dayside compression at 1640 UT





## Magnetic Field

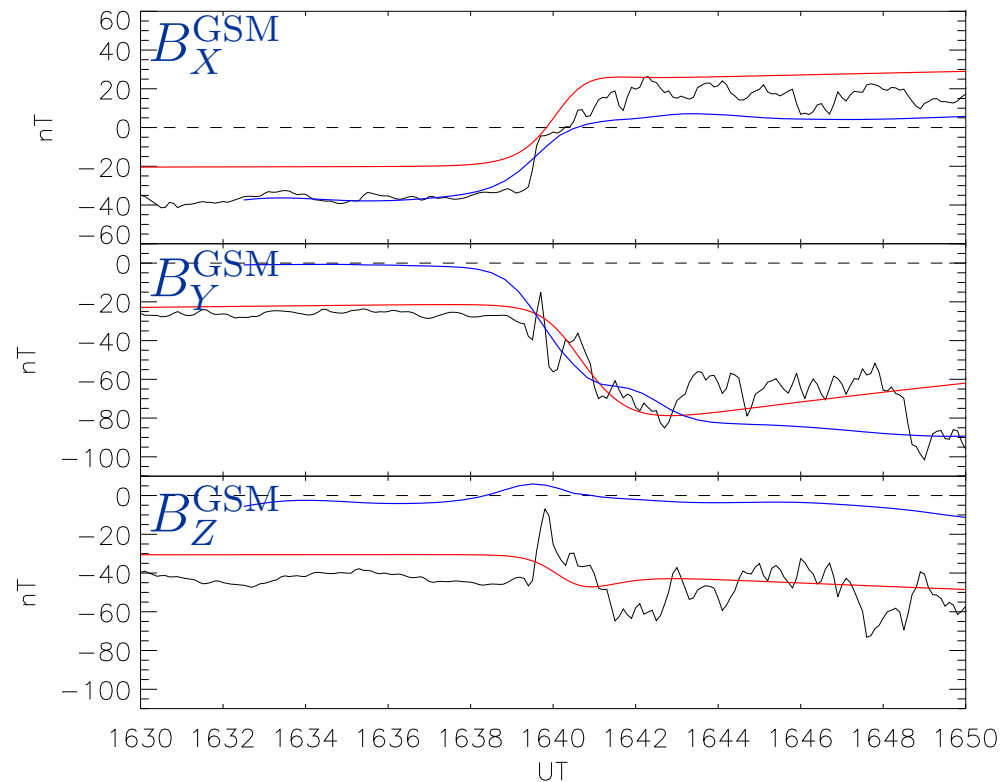
- GOES 8 located at noon close to the magnetic equator (see next slide)
- Measured field at GOES 8 (black)
- Modelled field (red)
- Magnetic field compression ( $B_Z^{\text{GSM}}$ )





## Magnetic Field (cont.)

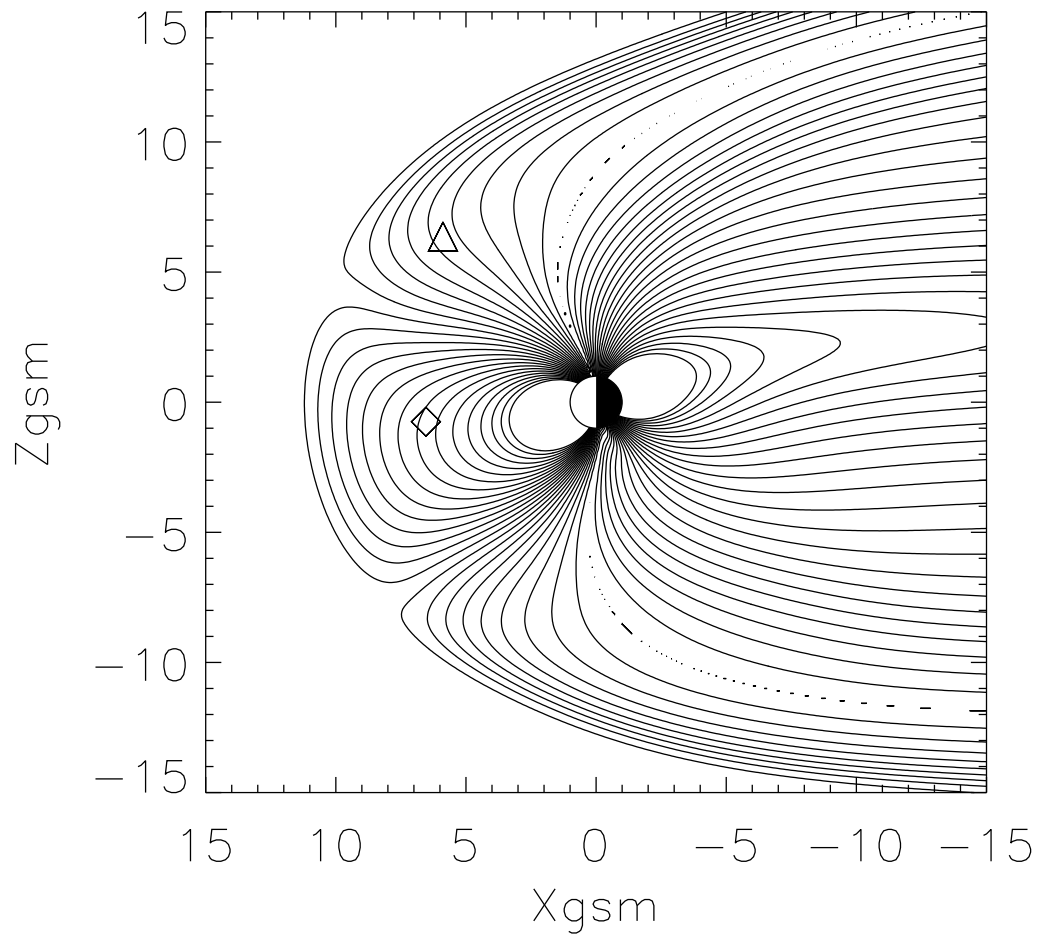
- Polar located at noon close to the northern cusp (see next slide)
- Measured field at Polar (black)
- MHD model from GU-MICS (blue)
- Magnetic field compression ( $B_X^{\text{GSM}}$ )
- Large deflection of  $B_Y^{\text{GSM}}$



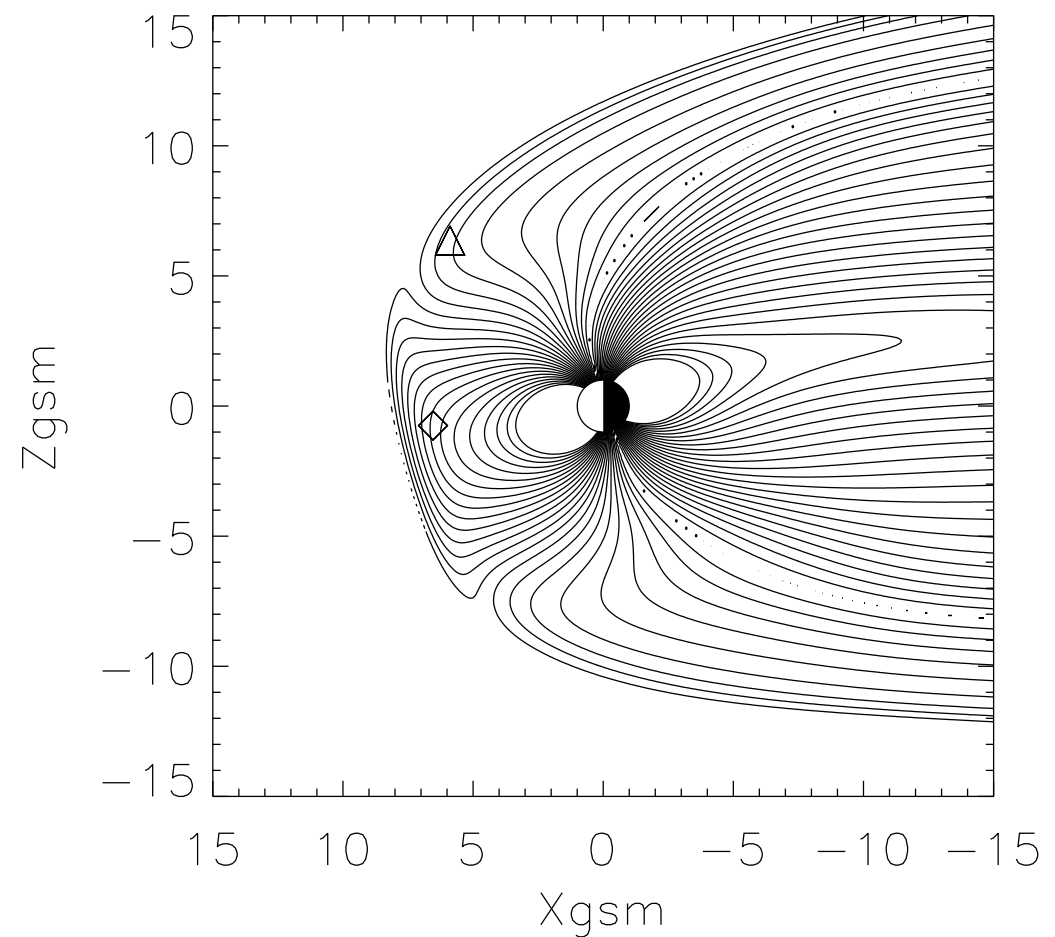


## Magnetic Field Configurations

1630 UT Polar( $\triangle$ ), GOES8 ( $\diamond$ )



1650 UT

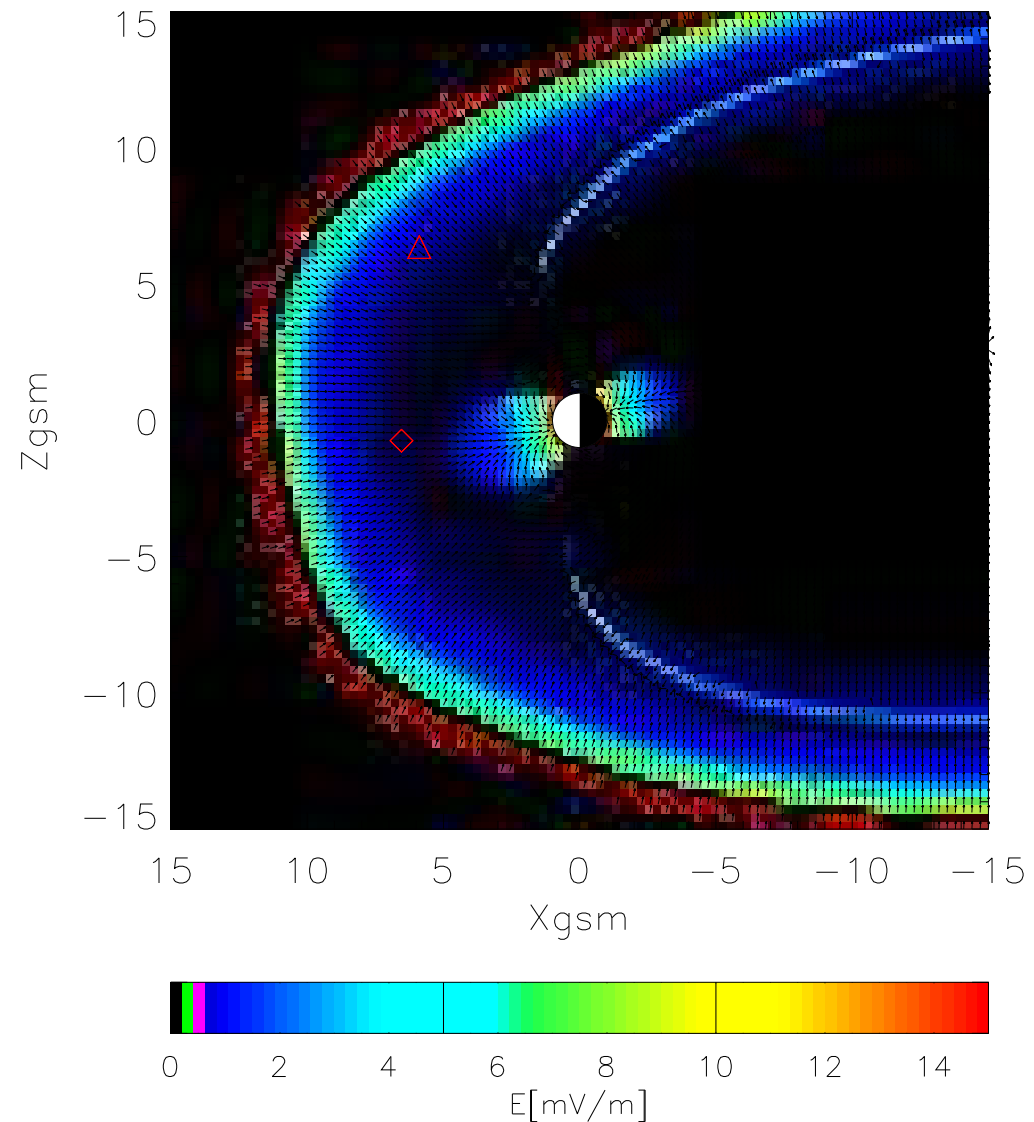






## Electric Field Configurations

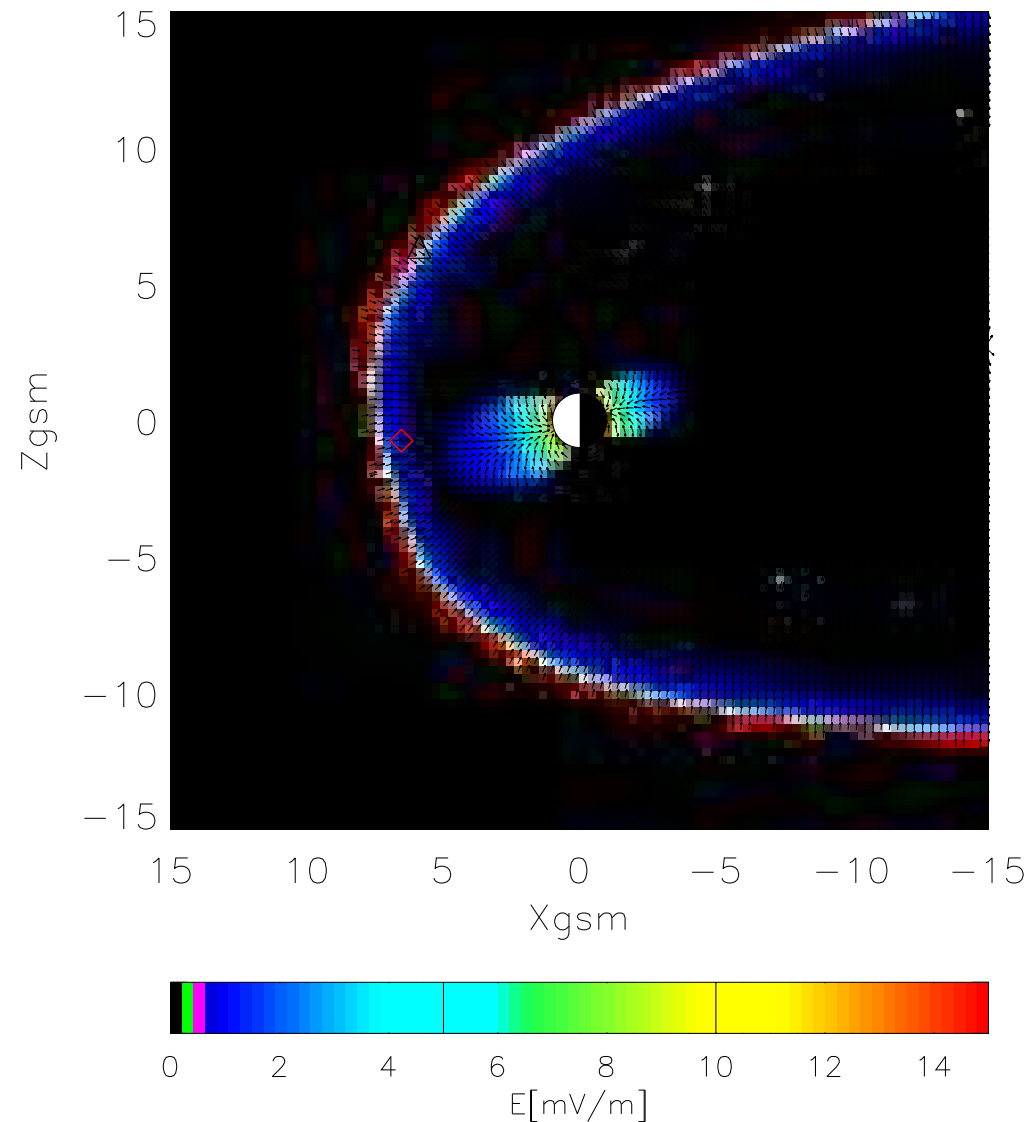
- Static electric field at 1630 UT
- Corresponds to the mapped ionospheric convection
- Determined by deformation to the initial dipole potential field
- $\triangle$  = Polar,  $\diamond$  = GOES 8





## Electric Field Configurations (cont.)

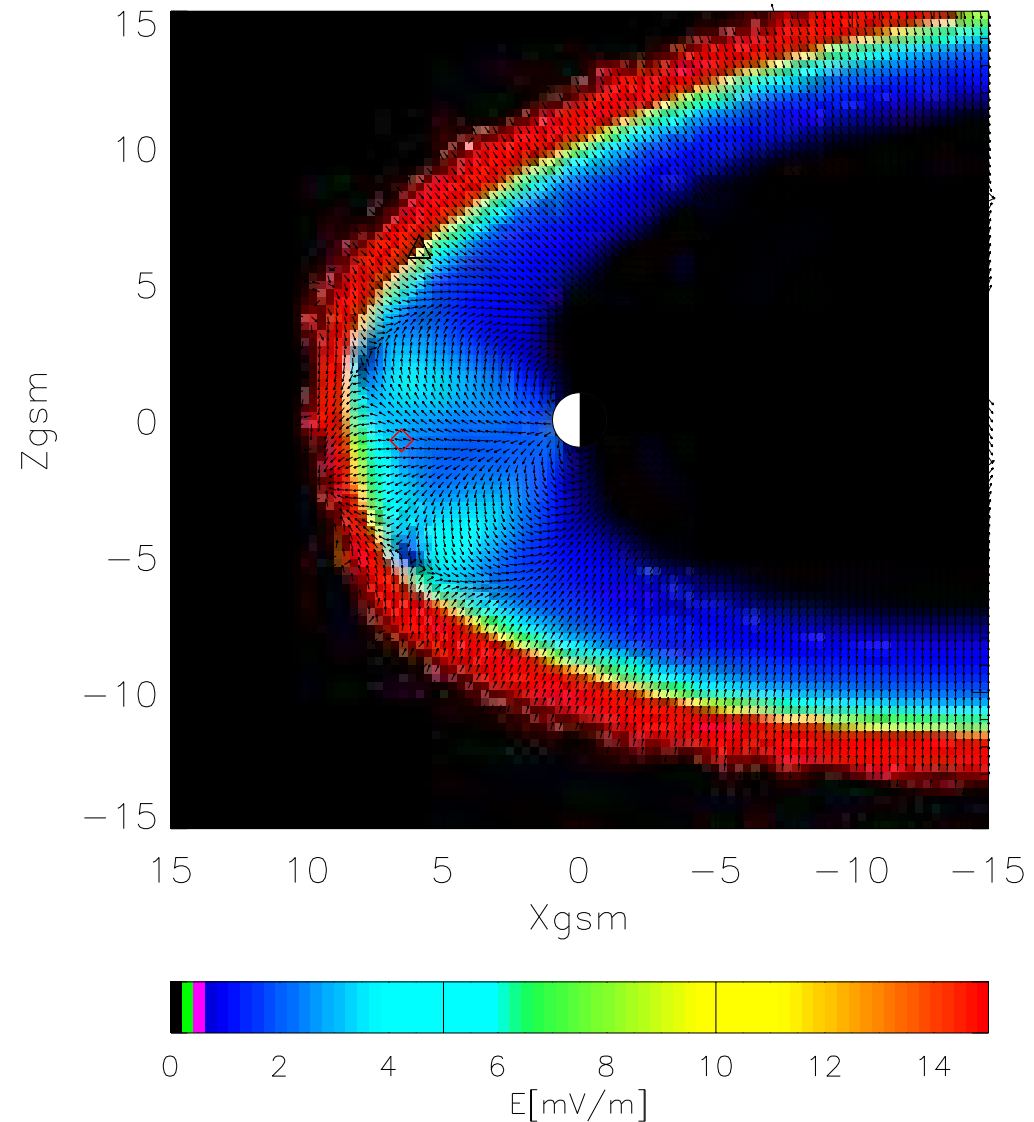
- Static electric field at 1650 UT
- $\triangle$  = Polar,  $\diamond$  = GOES 8
- Field close to the Earth is due corotation (plasmashpere)





## Electric Field Configurations (cont.)

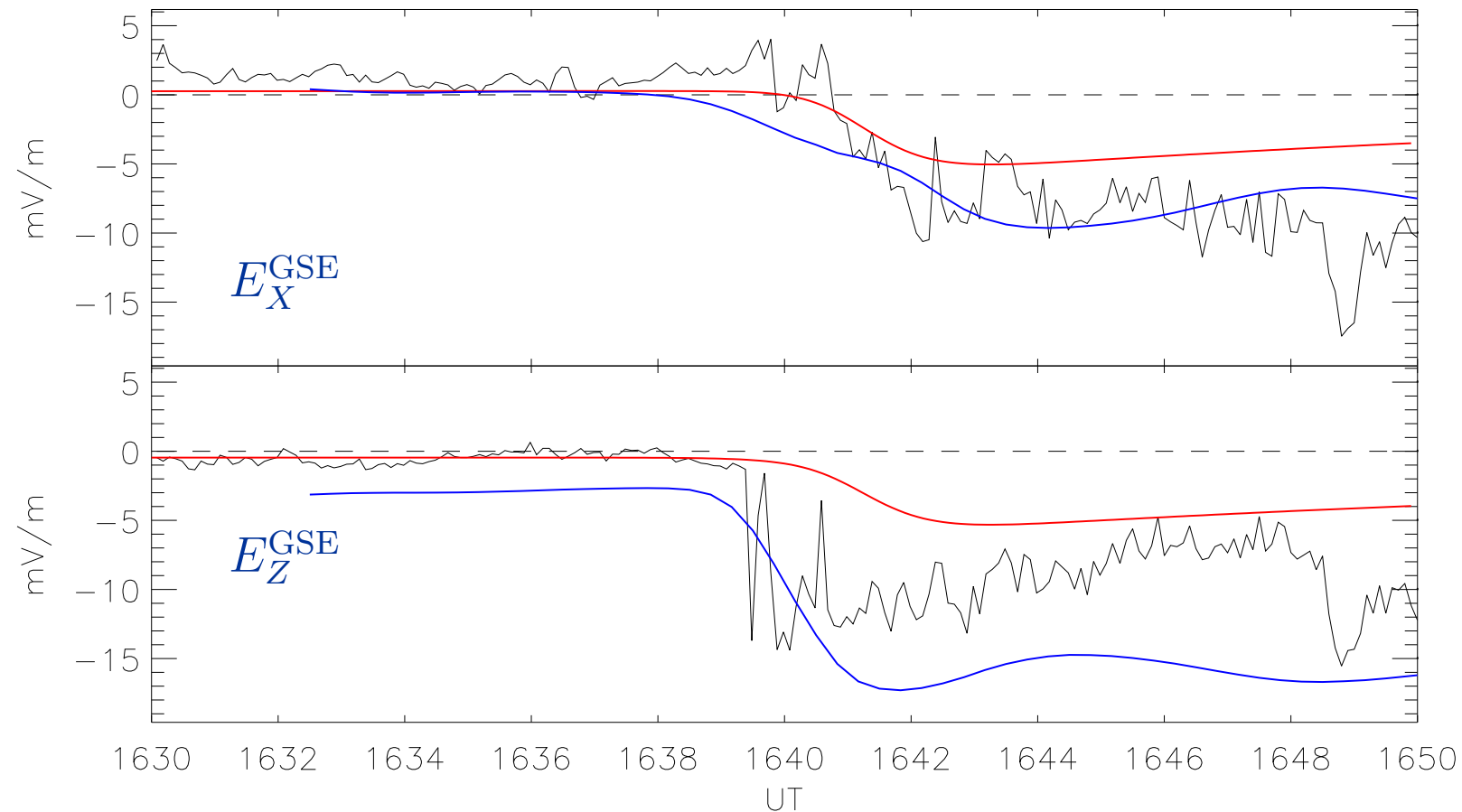
- Induced electric field at 1640 UT
- Corresponds to the field lines motion
- Due to time-evolution of  $B_Y^{\text{GSM}}$





## Electric Field at Polar

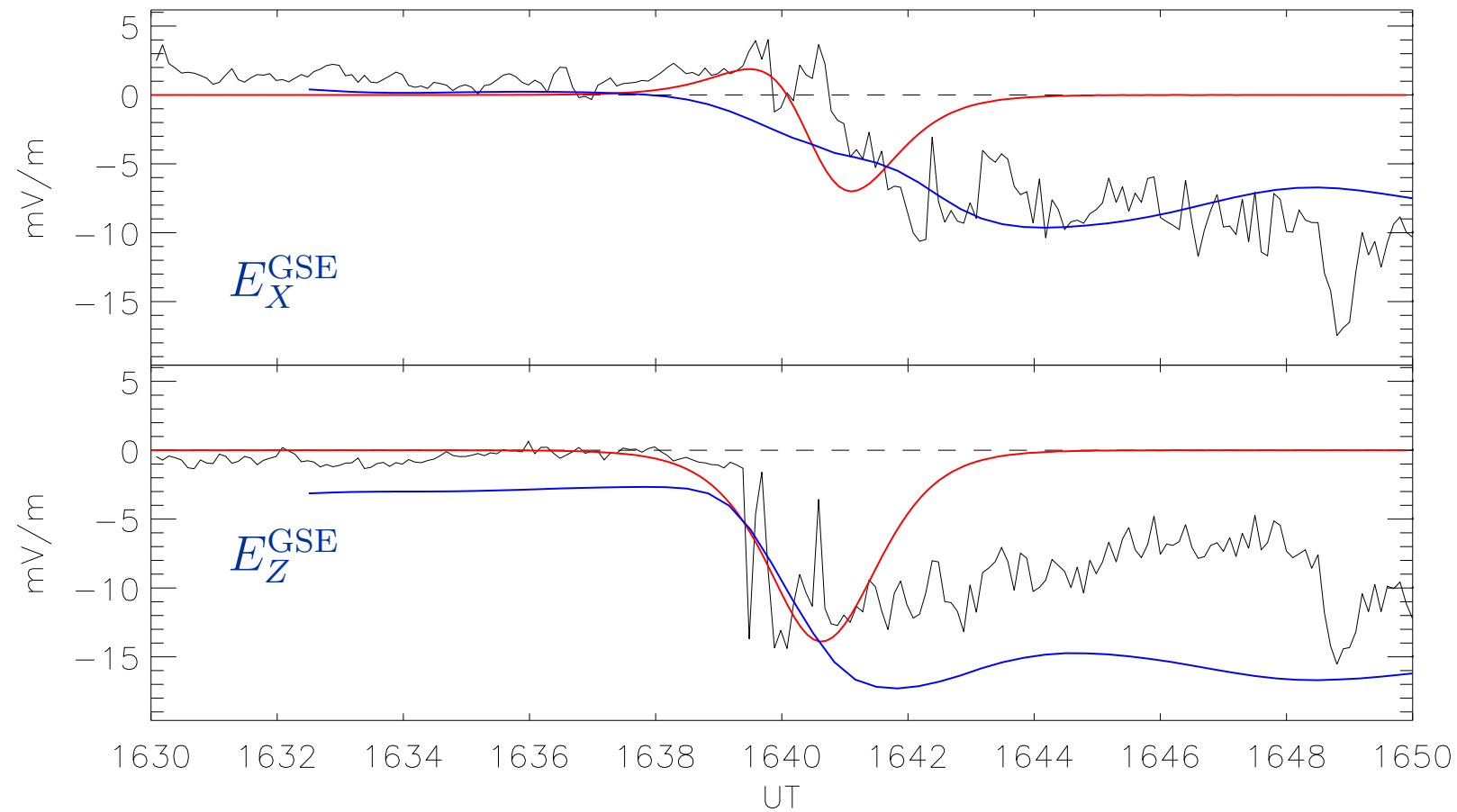
Convection electric field, Polar (black), modelled (red), GUMICS (blue)





## Electric Field at Polar (cont.)

Induced electric field, Polar (black), modelled (red), GUMICS (blue)







## Electric Field at Polar (cont.)

Total electric field, Polar (black), modelled (red), GUMICS (blue)

